Numerically Exact Many-Body Dynamics of Indistinguishable Particles

- User Manual -


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#### Abstract

This document is intended to guide any potential user to the installation and use of the MCTDH-X program package. MCTDH-X is a highly efficient numerical algorithm to solve the many-body interacting Schrödinger equation. The present manual describes the basic philosophy and structure of the program and the workflow of the MCTDH-X package. In particular, it explains how to operate and control the main program and the analysis routines from the respective input files.The output of the main program's numerical computations is visualized with bash scripts as images and video files with any desired plotting or data processing program. A description of the usage of predefined scripts is included in the present manual. Multiple computations, such as parameter scans can be automated with the provided "Monster" script. Finally, Mercurial, the version management system, is described and development and good programming guidelines are suggested.


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## 1 Installation

### 1.1 Prerequisites for the main program

The prerequisites to install the MCTDH-X program are Fortran compilers (Intel, PGI or GNU). Fast Fourier transforms need to be provided by either Intel MKL or FFTW libraries. An FFTW source is included in the package and will be compiled during the installation if FFTW is selected. LAPACK routines also have to provided as a library - either within the Intel MKL, or, if selected an included OpenBLAS source will be compiled and linked. A version of MPI, which provides an MPI Fortran compiler wrapper, such as mpif90, is also a prerequisite. For the successful compilation of the included Mercurial version management, python development headers have to be installed on the system. In Debian-based Linux distributions just install the packages gfortran, libopenmpi-dev, openmpi-common, openmpi-bin with your package/software management. In Ubuntu, for instance, you may do the following in terminal:
sudo apt-get install gfortran libopenmpi-dev openmpi-common openmpi-bin

### 1.2 Prerequisites for full functionality

To make the MCTDH-X software fully functional, the following packages are needed: python, libpython-dev, scons, gnuplot, mercurial, mplayer2, mencoder. With these packages, the integrated software management of the MCTDH-X package, and its visualization (movie) scripts should be fully functional. In Ubuntu, for instance, you may type the following in terminal to install the prerequisites:
sudo apt-get install python libpython-dev scons gnuplot mercurial mplayer2 mencoder

### 1.3 Running the installation

After the prerequisites are fulfilled, to install MCTDH-X one has to first unpack the program package (this unpacking can be skipped when you download from the repository with hg clone):

```
tar -xvf mctdhx.tgz
```

Subsequently, the program can be installed by running the interactive installation script:
./Install_MCTDHX.sh
The script offers several configuration options, like different platforms (Generic, mpif90, CrayXC40) and compilers (Intel, GNU, PGI,ftn). On a new platform, a modification of the Makefile or scons script might nevertheless be necessary. If the Generic build option is selected, the build is performed with gfortran and the included OpenBlas and FFTW sources. The mpif90 option will use either mpif90 or mpif90.openmpi for the build and the CrayXC40 option selects the Makefile.hornet for the Cray computer Hornet and uses the Cray Fortran compilers ftn. To get further information (if you are not at all familiar with GNU make and the build on your platform fails), contact and seek help at mctdhx@ultracold.org or in the forum http://ultracold.org/forum. The installation script creates a few aliases and appends to the \$PATH environment variable. It creates a file .mctdhxrc and .guantumrc and appends to the .bashrc to source the .mctdhxrc, such that the new commands become available. Available commands/aliases are collected in the following table 2 .

| Alias | What it does |
| :--- | :--- |
| cdm | takes you to the MCTDH-X installation directory |
| mcg.sh <X> | MCTDH-X code grep. Script which parses the MCTDH-X code and scripts <br> for the argument <X> |
| data_miner.sh | run an interactively configured analysis on computations in all subdirecto- <br> ries |
| MCTDHX | executes the the program (for interactive use, only. Most queueing systems <br> prevent aliases from working and the executables then have to be copied <br> and run in the current working directory) |
| MCTDHX_analysis | executes the the analysis program (for interactive use, only. Most queueing <br> systems prevent aliases from working and the executables have to be copied <br> and run in the current working directory) |
| libcp | copy the dynamic library libmctdhx.so to the current working directory |
| inpcp | copy the example inputs MCTDHX.inp, analysis.inp to the current direc- <br> tory |
| bincp | copies the executables of the main and analysis programs to the current <br> directory. |
| libmake | will recompile the dynamic library libmctdhx.so |
| x-make | will recompile the whole program |
| mctdhx_hg | alias to the supplied mercurial distribution (can be used for version man- <br> agement and to update the program)* |
| mctdhx_gnuplot | supplied gnuplot program to visualize data*. |
| mctdhx_mencoder | supplied mencoder program to make movies from gnuplot images*. |
| mctdhx_mplayer | supplied mplayer program to view the movies*. |

Table 2: Aliases and scripts provided with the MCTDH-X installation.
(aliases marked by a * are only created if respective included software is not present in the system)

## 2 An MCTDH-X Tutorial

A straightforward way of familiarizing youself with the usage of MCTDH-X, without further reading the remainder of this user manual, is to follow this step-by-step tutorial. The tutorial includes running the main program to compute a many-body eigenstate as well as to compute a timeevolution for bosons after changeing the potential. Subsequently, the analysis program is run to extract some quantities of interest that then are visualized with the included bash scripts to obtain videos of the computed dynamics.

### 2.1 Computing an Eigenstate by Relaxation

After the installation script has finished, you should make sure that the aliases and links are working. To do so, first type

```
source ~/.mctdhxrc
```


## ls \$MCTDHXDIR

. Now you should see the executables MCTDHX_<compiler>, MCTDHX_analysis_<compiler> and the library libmctdhx.so. If you don't, the installation did not terminate correctly and you should
check what went wrong (see the log files created in ./log/) and contact the developers in case you cannot find out or fix the error.

To get started, it's best to create a directory for the tutorial computations:

```
mkdir ~/MCTDH-X-Tutorial
cd ~/MCTDH-X-Tutorial
```

. To copy all necessary files to run the computation in this directory, we make use of the alias inpcp and libcp which copy the example inputs MCTDHX.inp, analysis.inp and the dynamic library libmctdhx. so to the current directory.

```
libcp
inpcp
```

ls
. The ls command should show a list including MCTDHX.inp, analysis.inp, and libmctdh.so. With these three files, we're able to run the main and analysis programs. The default MCTDHX.inp is configured to compute the eigenstate of $N=2$ bosons in a one-dimensional harmonic oscillator potential with unit frequency and with $M=4$ orbitals. To make things a little more interesting, let's change the the number of bosons to $N=50$ by editing MCTDHX.inp and setting Npar $=50$ in the System_Parameters namelist, i.e.,

```
gedit ./MCTDHX.inp
```

. Since we're going to compute some dynamics later, we use the opportunity to also enlarge the number of grid points and the grid extension in the DVR_Parameters namelist (a bit further down in the MCTDHX. inp file). We set:

```
Npar = 50
NDVR_X = 256
x_initial = -12.d0
x_final = 12.d0
```

. With these adjustments made, we can run the relaxation to the groundstate of the harmonic oscillator potential by typing

## MCTDHX

. This should show an output similar to the following figure 1. This computation is going to take a minute so sit back and relax or just get a coffee, until it finished. The energy you should see on screen in the final propagation step should be identical to 259.14031953 . To visualize the output, let's use gnuplot:

```
mctdhx_gnuplot
plot "20.0000000orbs.dat" u 1:8, "" u 1:(sqrt($24**2)), "" u 1:(sqrt($22**2))
```

. The plot command visualizes the density $\rho(x)$ (column 8) and the first two natural orbitals $\phi_{1}^{(N O)}(x)$ and $\phi_{2}^{(N O)}(x)$ (column 24 and 22, respectively) in the last ASCII output file of the relaxation 20.0000000 orbs.dat. For a full explanation of the structure of this file, see table 14 . In figure 2, you can see a screenshot of what your gnuplot output should look alike. To find out the occupations of the natural orbitals, let's have a look at the last line of the NO_PR. out file:

MCTDH-X


Figure 1: Output of an MCTDH-X computation in the shell.


Figure 2: Visualization of MCTDH-X relaxation using gnuplot. The density $\rho(x)$ and the first two natural orbitals $\phi_{1}^{(N O)}(x)$ and $\phi_{2}^{(N O)}(x)$ are shown as red, green and blue line, respectively.

## tail -n 1 NO_PR.out

We get
which is the time $t$, the natural occupations $\rho_{M}^{(N O)}(t), \rho_{M-1}^{(N O)}(t), \ldots, \rho_{1}^{(N O)}(t)$ and the energy of the system $E(t)$, in the final column (see also table 11 for the structure of the NO_PR. out file). We conclude that despite the strong interactions of $\lambda_{0}=1.0$, our eigenstate of the $N=50$ bosons in the
harmonic confinement is close to condensed, since $96 \%$ of the bosons sit in the lowest single-particle state. In the single-well, this absence of fragmentation in the groundstate is anticipated. There is, however, numerous examples on the emergence of fragmentation in the dynamics of ultracold bosonic systems. To see the occurence of fragmentation, it's therefore instructive to change the potential in which our eigenstate was computed and trigger some dynamics.

### 2.2 Computing the Time-Evolution of a System

To propagate a given initial state in time, one has to change only a few parameters in the MCTDHX.inp file. To start, it's best to create a subdirectory for the propagation inside the MCTDH-X-Tutorial directory which contains the relaxation:

```
cd ~/MCTDH-X-Tutorial
mkdir propagation-double-well
```

. As indicated by the name, we're going to propagate the groundstate of the harmonic potential, which we computed in the previous subsection, in a double well potential. First, we copy the necessary files to the newly created directory:

```
cp MCTDHX.inp libmctdhx.so analysis.inp PSI_bin CIc_bin ./propagation-double-well/
cd propagation-double-well
```

Subsequently, the input file needs to be adapted with a text-editor, i.e.,

```
gedit MCTDHX.inp
```

. To make the program propagate the initial state in the binary data files PSI_bin and CIc_bin, the following parameters have to be adapted:

```
Job_Prefactor=(0.d0,-1.d0)
GUESS='BINR'
Binary_Start_Time=20.0d0
```

as also explained in the inlined documentation in the MCTDHX.inp file and in table 3, Job_Prefactor $=(0 . d 0,-1 . d 0)$ triggers the program to do a forward time propagation, GUESS $=$ 'BINR' will make it read the initial state from the binary files CIc_bin and PSI_bin, and Binary_Start_Time $=20 . \mathrm{d} 0$ chooses the time at which the binary files CIc_bin and PSI_bin are read for the initial state. Now, in order to define the parameters specifying the propagation in the double well, we change the following variables in the MCTDHX.inp :

```
Integration_Stepsize=0.0001d0
whichpot="h+d"
parameter1=2.d0
parameter2=8.d0
parameter3=1.d0
```

. Integration_Stepsize $=0.0001 \mathrm{~d} 0$ specifies the initial time-step (compare table 3). The which_pot $=$ ' ' $h+d$ '' variable selects a potential which is the sum of a displaced harmonic potential and a discplaced Gaussian barrier in its center. To define the potential parameter1=2.d0 is the displacement, parameter $2=8 . d 0$ specifies the height of the barrier in the center of the parabola and parameter3=1.d0 gives the width of this barrier. Finally, we select the appropriate integrator to propagate the coeffcients' equations of motion:


Figure 3: Checking the status of an MCTDH-X computation with mctdhx_status.sh.
The left panel displays the time-evolution of fragmentation, i.e., $1-\sum_{i=2}^{M} \rho_{i}^{(N O)}(t)$, in the system and the right panel shows a snapshot of the current density.

```
Coefficients_Integrator='MCS'
```

. Now the input is adjusted to start the propagation and we type

## MCTDHX

to run it and wait until it finished. After it has finished, we can get a quick visualization of the final state by typing

```
mctdhx_status.sh
gnome-open status.png
```

. The script mctdhx_status.sh plots the time-evolution of the fragmentation in a computation as well as a snapshot of the current density. The picture you should see is displayed in Figure 3. Finally, we could invoke several different bash-scripts to visualize the computed time-evolution as movies (cf. table 9) or simply run all available scripts with the visualization master script, like so:

```
vms.sh all $PWD -7.0 9.0
```

(cf. also section 4.3). vms.sh automatically runs the analysis program on the data in the present directory and creates videos of the density as well as the correlation functions of the system. For reference, these are also available on the website at http://ultracold.org/documentation.

## 3 Program Structure

The MCTDH-X package contains a main program to perform the actual numerical task and an analysis program that is used to compute the desired quantities of analysis from the many-body wavefunction. For the purpose of visualization, bash scripts that generate .mpg or .avi video files by running gnuplot and mencoder are provided. A source tarball of gnuplot and mencoder are provided with MCTDH-X in the External_Software subdirectory and installed with the installation script,
as mentioned above. To automate and simplify the use of the main program, the analysis program as well as the visualization bash scripts, the program package also contains a graphical user interface (Guantum), as well as scripts that automate series of computations (MonsterScript.sh) and a script for the automation of data processing (data_miner.sh).

### 3.1 Main Program

The main program can be run by typing MCTDHX (in wrappers or in runscripts the MCTDHX alias might not work, and one needs to use MCTDHX_intel, MCTDHX_pgf or MCTDHX_gcc). Although, in the case of a computationally intense task, it is a lot faster to use the shared memory and distributed memory parallelization of the program and run it with one the various instances of MPI launchers (such as mpirun,mpiexec, mpiexec.hydra, aprun, ...). Examples for running MCTDHX in parallel can be found in the example PBS scripts directory PBS_Scripts. Since the configuration of such an hybridly parallel job is complicated, it is easier to just use the MonsterScript.sh script that will automate whole series of hybridly parallel computations. They can be found in the ./Computation_Scripts directory. If a manual configuration of a task is needed or desired, this is done by adapting one of the PBS runscript examples. Depending on the hardware architecture, the most efficient way is usually to run MCTDH-X with at least as many MPI processes as there are orbitals. The OpenMP shared memory parallelization takes care of efficiently performing the computational task inside of each MPI process. The program's structure is entirely modular, i.e., all subroutines are collected in Fortran modules. To inspect the program structure, please consult the html documentation by opening the index.html file in the documentation/html subdirectory. In this html documentation all important variables are explained and for each routine call- and caller-graphs are given.

### 3.2 Analysis Program

The analysis program can be run by typing MCTDHX_analysis. The analysis program is serial, i.e., no shared or distributed memory parallelization is used so far. The program is also entirely modular and the structure, call- and caller graphs, as well as a documentation of important variables is available in the same html documentation as for the main program (see file documentation/html/index.html).

### 3.3 Scripts

In the subdirectory ./bin/Scripts there is the MonsterScript.sh script that can configure whole series of computations. The input file ParameterScan.inp is in the ./Input_Examples directory and contains an in-line documentation. MonsterScript.sh allows a scan of 5 user-defined parameters for the relaxation, and 5 user-defined parameters for the subsequent propagations. A parameter scan for potential, DVR, particle or orbital number, and magnitude and/or width of the interaction can be configured, for instance. The configuration of MonsterScript.sh is detailed in table 10. The directory ./Computation_Scripts contains the visualization master script, vms.sh, that automatically runs the MCTDHX analysis program and visualizes a given computation. The possible different visualization videos and plots are realized as bash scripts which can be found in the ./Visualization_Scripts directory.

## 4 Input File driven Usage MCTDH-X

The basic workflow of a numerical solution of the time-dependent many-boson Schrödinger equation, $\hat{H}|\Psi\rangle=i \partial_{t}|\Psi\rangle$ (TDSE), with MCTDH-X is the following:

1. Define the Hamiltonian $\hat{H}$.

Modify Get_InterParticle_Potential.F and Get_1bodyPotential.F.
2. Define the initial state $|\Psi\rangle$.

Modify Get_Initial_Coefficients.F and Get_Initial_Orbitals.F.
3. Solve the TDSE by propagating in real or imaginary time.
libmake, then libcp, (maybe inpcp, modify MCTDHX.inp), finally MCTDHX.
4. Analyze and visualize the solution $|\Psi\rangle$.
(maybe inpcp, modify analysis.inp), run MCTDHX_analysis, process ASCII data with gnuplot.

Steps (1) through to (3) are done with the main program and step (4) is done with the analysis program and movie bash scripts. There is several ways of automatization of the above 4 -step scheme. As mentioned above, there is a scripts called MonsterScript.sh in the Computation_Scripts directory that allows a fully integrated and automated processing of whole series of computations. Finally, the data processing to videos is automated in the visualization master script vms.sh.

### 4.1 Defining the Hamiltonian

The Hamiltonians that can be treated by the MCTDH-X package in its current implementation are those with maximally two-body operators. In general such a Hamiltonian hence contains the kinetic energy $\hat{T}$ and the external potential $\hat{V}$ and an interparticle interaction $\hat{W}$ :

$$
\begin{equation*}
\hat{H}=\sum_{i=1}^{N}\left(\hat{T}_{i}+\hat{V}_{i}\right)+\sum_{i<j=1}^{N} \hat{W}_{i j} \tag{1}
\end{equation*}
$$

To fully define this Hamiltonian in dimensionless units [see Phys. Rev. A 77, 033613 (2008)] one has to define the one-body potential $\hat{V}$, and the two-body interaction $\hat{W}$. It has to be stressed, that both operators can be in principle time-dependent.

## Specifying the one-body potential $\hat{V}$ :

The one-body potential $\hat{V}$ is specified in the file source/ini_guess_pot/Get_1bodyPotential.F. This file contains a Fortran subroutine with a case selection for the dimensionality of the treated problem. The potentials defined in this routine can be selected and modified with the input parameters whichpot, parameter1, parameter2, ..., parameter30 in MCTDHX.inp. Appendix A collects the predefined potentials and parameters. If a custom potential is desired, it has to be implementated in the file source/ini_guess_pot/Get_1bodyPotential.F (look there for the loops which are enclosed by the if- exceptions for whichpot .eq. custom1D/custom2D/custom3D, respectively).

## Specifying the two-body potential $\hat{W}$ :

The two-body interaction $\hat{W}$ is specified in the file Get_InterParticle_Potential.F in the directory source/ini_guess_pot/. This file contains a single Fortran subroutine with a case selection on the type of interparticle interaction and the evaluation of its action (Interaction_Type is the corresponding input variable). The standard types are a contact interaction potential (Interaction_Type $=0,6$ ) or short range Gaussian interaction (Interaction_Type=1,2,3,4,5). For these cases of a short-range interparticle interaction, the width of it can be adjusted with the input variable Interaction_Width in MCTDHX.inp. Generally, a non-contact interaction requires Interaction_Type > 0 . Interaction_type $=5$ is a time-dependent non-contact interaction and Interaction_type $=6$ allows the simulation of contact interactions with a time-dependent interaction strength. The different evaluation types that Interaction_Type selects and the necessary properties of the interaction potentials are also specified in the html documentation and the example input in the directory Input Examples. If a custom interaction potential is desired, this has to be implemented in the Get_InterParticle_Potential.F routine. The table in Appendix B shows the predefined interparticle interaction potentials as well as how to configure them with the input file's parameters.

## Defining the initial state $|\Psi\rangle$ :

The MCTDH-X wavefunction is a multiconfigurational expansion with time-dependent coefficients $C_{\vec{n}}(t)$ and time-dependent configurations $|\vec{n} ; t\rangle$. In order to fully define the wavefunction,

$$
\begin{equation*}
|\Psi(t)\rangle=\sum_{\{\vec{n}\}} C_{\vec{n}}(t)|\vec{n} ; t\rangle \tag{2}
\end{equation*}
$$

one has to define all coefficients $C_{\vec{n}}$ and all configurations $|\vec{n} ; t\rangle$. In the case of a fully user-defined guess, i.e., GUESS='HAND' in the input file, the wavefunction is supplied via Fortran routines. The coefficients are defined in the file source/ini_guess_pot/Get_Initial_Coefficients.F and the orbitals from which the configurations are built are defined in the file Get_Initial_Orbitals.F in the directory source/ini_guess_pot/. Both contain a single Fortran subroutine, that assigns the corresponding array (for a documentation of the routines, please see the html documentation). For a relaxation, i.e., the calculation of a certain eigenstate of the many-body Hamiltonian $\hat{H}$, usually it is a good choice to start from a Gross-Pitaevskii, i.e., single-configurational state with $|\Psi\rangle=|N, 0,0, \ldots\rangle$. For a propagation there is also the possibility to restart the computation from a previous one by specifying GUESS= 'BINR' in the input file.

### 4.2 Running the computation and analysis

In this subsection, it is specified how to run the MCTDH-X program and analysis manually, i.e., by configuring the MCTDHX.inp and analysis.inp input files and running the respective programs. To read how to automate this process using the MonsterScript.sh and vms.sh scripts, please consult the next subsection 4.3.

## The input of the main program

To run the MCTDH-X program, it is best to have each calculation done in a separate directory. The program needs two files to perform the computation. First, the dynamic library, libmctdhx.so, and second, the input file with computational/numerical parameters MCTDHX.inp. The library contains the above four subroutines that define the Hamiltonian and the initial state (Get_Initial_Coefficients.F, Get_Initial_Orbitals.F, Get_InterParticle_Potential.F, Get_1bodyPotential.F). If any of these files was modified, the library has to be recompiled and copied to the computation's directory! Recompiling the library libmctdhx. so is conveniently achieved by issueing the command:
libmake
After compilation, the library can be copied to the working directory by typing

## libcp

Finally, the input files should be copied to the working directory. This can be done by typing

```
inpcp .
```

The two files MCTDHX.inp and analysis.inp are now inside the working directory. The details of the input variables in those files can be taken from the in-line documentation in the input file or the html code documentation. The MCTDHX.inp file basically contains variables like the particle number, the number of orbitals, the (dimensionless) interaction strength, number of primitive basis functions, their type, details of the integration (integrator,stepsize,accuracy,order,...), and many more. See table 3 for all currently available parameters in the main program and their meaning.

| System Parameters Namelist |  |  |
| :---: | :---: | :---: |
| Parameter | Meaning | Options |
| JOB_TYPE | Select problem, i.e. MCTDH for bosons, MCTDH for fermions or TDCI for bosons. | Character, 'BOS', 'FER', 'FCI'; Default 'BOS' |
| Morb | Select number of time-dependent, variationally optimized basis functions | Integer, no default. |
| Npar | Select number of structureless particles | Integer, no default. |
| xlambda_0 | Adjust prefactor of two-body potential in the Hamiltonian | Real, no default. |
| mass | Mass of the particles | Real, default 1.d0 |
| Job_Prefactor | Select which direction to propagate the equations of motion in time. | Complex, $(0.0 \mathrm{~d} 0,-1.0 \mathrm{~d} 0)$ <br> Forward propagation; <br> $(0.0 \mathrm{~d} 0,+1.0 \mathrm{~d} 0)$ Backward <br> propagation, $(-1.0 \mathrm{~d} 0,0.0 \mathrm{~d} 0)$ <br> (Improved) Relaxation; <br> Default (-1.d0,0.d0).  |
| NProjections | Number of times that the projection operator $\hat{\mathcal{P}}$ is applied to the right-hand side of the orbital equations of motion | Integer, default 2. |


| GUESS | Specifying if the initial guess is defined in .dat files, the routines Get_Initial_Orbitals.F, Get_Initial_Coefficients.F, or in the binary files CiC_bin and Psi_bin | Character, 'HAND' means the Get_Initial. . . routines are used, 'DATA' means .dat files will be read in, 'BINR' means *_bin files will be read in, default 'HAND'. |
| :---: | :---: | :---: |
| Diagonalize_OneBodyh | Select to use eigenfunctions of the one-body potential in Get_1bodyPotential.F | Logical, .T. or .F., .T. only non-FFT ${ }^{1}$ DVR迆, i.e., DVR_X/Y/Z $\neq 4$, default is .F. |
| Binary_Start_Time | Define point in time at which the wavefunction is read from binary files in the case of GUESS='BINR' | Real, no default. |
| Restart_State | If restarting from a Block Davidson computation, this selects which state in the block is used as the initial value. | Integer, default 1 |
| Restart_Orbital_FileName | Define filename of orbitals to read if GUESS= ' $D A T A$ ' | Character, no default. |
| Restart_Coefficients_FileName | Define filename of coefficients to read if GUESS= 'DATA' | Character, no default. |
| Vortex_Seeding | In propagations: Multiply initial orbitals with phase/density profile? | Logical, default .F. |
| Vortex_Imprint | In relaxations: Apply a projection operator to a certain orbital density profile? | Logical, default .F. |
| Profile | If Vortex_Imprint is true, this will select the respective shape of the imprinted (select 'tanh','poly','phase', or 'phase-x' or define custom profile in Get_Vortex_Profile.F in /source/ini_guess_pot/) | Character, default 'tanh' |
| Fixed_LZ <br> OrbLz | Multiply fixed phase profile on the orbital in relaxations <br> How many times $2 \pi$ the phase is going to jump if Fixed_LZ is true. If any value is -666 , the respective orbital is unchanged. | Logical, default.F. Integer, $0,0,0,0,0,0,0,0,0,0 \quad$ default |
| DVR Namelist |  |  |
| DIM_MCTDH | Specifying the dimensionality of the problem | Integer, 1,2 or 3, default 1 . |
| NDVR_X | Specifying the number of DVR functions in X dimension | Integer, default 256. |

[^0]| NDVR_Y | Specifying the number of DVR functions in Y dimension | Integer, default 1. |
| :---: | :---: | :---: |
| NDVR_Z | Specifying the number of DVR functions in Z dimension | Integer, default 1. |
| DVR_X | Which DVR will be used in X direction | Integer, 1 means harmonic oscillator DVR, 3 means sine DVR, 4 means FFT DVR and 5 exponential DVR, default 4. |
| DVR_Y | Which DVR will be used in Y direction | Integer, 1 means harmonic oscillator DVR, 3 means sine DVR, 4 means FFT DVR and 5 exponential DVR, default 4. |
| DVR_Z | Which DVR will be used in Z direction | Integer, 1 means harmonic oscillator DVR, 3 means sine DVR, 4 means FFT DVR and 5 exponential DVR, default 4. |
| x_initial | Where the spatial grid in X dimension starts | Real, default -8.0. |
| x_final | Where the spatial grid in x dimension stops | Real, default 8.0. |
| y_initial | Where the spatial grid in Y dimension starts | Real, default -8.0. |
| y_final | Where the spatial grid in Y dimension stops | Real, default 8.0. |
| z_initial | Where the spatial grid in Z dimension starts | Real, default -8.0. |
| z_final | Where the spatial grid in Z dimension stops | Real, default 8.0. |
| Integration Namelist |  |  |
| Time_Begin | Time at which the simulation shall start | Real, default 0.0. |
| Time_Final | Time at which the simulation shall stop | Real, default 20.0. |
| Time_Max | Maximal time | Real, default 1.d99. |
| Output_TimeStep | Times at which orbital output shall be written. | Real, default 0.1. |
| Output_Coefficients | Multiple of Output_TimeStep at which coefficient output shall be written | Real, default 1 |
| Integration_Stepsize | Stepsize of the integration scheme (for relaxation this is fixed, for propagation this is adaptive) | Real, default 0.01 |
| Error_Tolerance | Error tolerance for the integration scheme | Real, default 1.d - 9 |


| Minimal_Occupation | Minimal occupation for not considering an eigenvalue as 0 in the inversion of the density matrix elements | Real, default 1.d - 12 |
| :---: | :---: | :---: |
| Minimal_Krylov | Minimal size of Krylov basis for coefficients' integration | Integer, default 4 |
| Maximal_Krylov | Maximal size of Krylov basis for coefficients' integration | Integer, default 20 |
| Orbital_Integrator | Integrator for the orbital equations of motion | Character, 'ABM' or 'OMPABM' means that (OpenMP parallelized) Adams Bashforth Moulton predictor corrector integrator is used, 'BS' means Bullirsch-Stör, 'RK' means Runge-Kutta, and 'STIFF' means ZVODE specialized to stiff equations is used; default 'OMPABM'. |
| Orbital_Integrator_Order | Order of the integration of the orbitals' integrator, choice depends on the Integrator | Integer, for 'RK' it is 5 or 8, for 'ABM'/'OMPABM' it is 2 to 8 , for ' BS ' its 2 to 16 , and for 'STIFF' it is 1 or 2 , default 7. |
| Orbital_Integrator_MaximalSterRestricts the maximum stepsize in the orbital equations' integration |  | Real, default 0.01. |
| Write_ASCII | Specifying wether ASCII files are output during the computation | Logical, default .T.. |
| Error_Rescale_TD | Specifying a scale for the Error_Tolerance parameter for timedependent one-body potentials. | Real, default 1.0. |
| LZ | Trigger one-body angular momentum operator in the Hamiltonian | Logical, .T. or .F., default .F. |
| OMEGAZ | Prefactor of one-body angular momentum operator. | Real, default 0.0. |
| STATE | Which eigenstate of the Hamiltonian to compute. | Integer, 1 means the groundstate, default 1. Stable for Coefficients_Integrator $={ }^{\prime}$ DSL' or 'DAV'. |


| Coefficients_Integrator | Which integrator to use for the coeffi- <br> cients equations of motion. | Character, 'MCS' means <br> MCTDH SIL routine, 'DSL' <br> means MCTDH SIL diag- |
| :--- | :--- | :--- |
|  |  | onalization routine, 'DAV' <br> means Davidson diagonal- <br> ization, 'BDV' means block <br> Davidson diagonalization; no <br> default. |
| BlockSize | If Block Davidson (BDV) is used as <br> integrator for the coefficients then this <br> selects the number of coefficient vec- <br> tors in the block |  |
| Integer, default 4 |  |  |


| Interaction_Type | How the Local interaction potentials $\hat{W}_{s l}$ and their action is evaluated | Integer, 0 means $\delta$-like contact interaction potential, 1 to 4 will use the routine Get_InterParticle_Potential.F to generate the interaction potential. 1 means the potential is separable and hence one potential vector is allocated for each spatial dimension; 2 means the potential depends on the distance of the particles only, hence, a tridiagonal representation is used (only for aequidistant DVRs 3,4 and 5); 3 means full interaction matrix will be stored (very large array!!!), 4 interaction matrix evaluated with successive FFT (IMEST); 5 means time-dependent interaction with IMEST; 6 means time-dependent contact interaction; 7 means both contact and non-zero ranged interactions with IMEST; default 0 . |
| :---: | :---: | :---: |

Table 3: MCTDH-X main program input file parameters.

## Special parameters to treat multi-level atoms and spinors

In order to treat particles that have internal structure, like multileveled atoms or spinor particles, a set of special input variables has been introduced to the above System Parameters Namelist. These define the number of levels, the presence of a conical intersection and whether the interparticle interaction contains a spin-dependent part or not. In principle, a conical intersection amounts to off-diagonal terms in the one-body Hamiltonian $\hat{h}_{i}^{(C I)}$ which couple the different levels whereas a spin-dependent interparticle interaction means that there is two-particle interactions $\hat{W}_{\text {spin }}$ that can change the spin of particles. These terms read as follows:

$$
\begin{align*}
\hat{h}_{i}^{(C I), k k} & =\left(\hat{T}_{i}+V_{k k}\left(\hat{\vec{r}}_{i}\right)\right) \otimes \mathbf{1}_{\vec{r}} ; \quad \hat{h}_{i}^{(C I), k j}=V_{k j}\left(\vec{r}_{i}\right) \hat{\pi}_{j k}  \tag{3}\\
\hat{W}_{\text {spin }} & =\sum_{\nu=x, y, z}\left(\mathbf{S}^{\nu} \otimes \mathbf{1}_{\vec{r}}\right) \hat{W}\left(\vec{r}, \vec{r}^{\prime} ; t\right)\left(\mathbf{S}^{\nu} \otimes \mathbf{1}_{\vec{r}^{\prime}}\right) . \tag{4}
\end{align*}
$$

Here, the operator $\hat{\pi}_{j k}$ was defined, which makes level $j$ appear in the coordinate space of level $k$ for the spinor orbital on its right. Furthermore, the representation $\mathbf{S}^{\nu}$ was introduced for the general spin operators in $\nu=x, y, z$ direction. For problems including spin-orbit interactions the folloing Hamiltonian $\hat{h}_{S O}$ is added to the one-body Hamiltonian:

$$
\begin{equation*}
\hat{h}_{S O}=\gamma\left[\alpha\left(\hat{p}_{x} \mathbf{S}^{y}-\hat{p}_{y} \mathbf{S}^{x}\right)+\beta\left(\hat{p}_{x} \mathbf{S}^{y}+\hat{p}_{y} \mathbf{S}^{x}\right)\right] . \tag{5}
\end{equation*}
$$

The parameters $\alpha, \beta, \gamma$ are the prefactor of the Rashba spin-orbit term, Dresselhaus spin-orbit term and the overall strength of the spin-orbit coupling (see also input parameters below). Generally, when the above terms are present in the Hamiltonian, a transfer of population between the different levels of the treated particles is allowed. The input variables necessary to control the program through the MCTDHX.inp file in the case of multileveled or spinor particles are specified in the following table 4.2 .

| System Parameters Namelist |  |  |
| :--- | :--- | :--- |
| Parameter | Meaning | Options |
| NLevel | How many levels do the considered <br> particles have? | Integer, default 1. |
| Multi_level | Do the atoms have internal struc- <br> ture? | Logical, default .F. |
| Conical_Intersection | Does the one-body Hamiltonian con- <br> tain terms $V_{j k}(\vec{r})$ that couple differ- <br> ent internal states? | Logical, default .F.. If <br> set to .T., the potential <br> VTRAP_EXT which is defined in <br> Get_1bodyPotential.F con- <br> tains one additional vector <br> that stores $V_{j k}$ |
| InterLevel_InterParticle | Does the interparticle interaction <br> couple different internal states di- <br> rectly (attention: this is for mul- <br> tileveled atoms which are NOT <br> spinors | Logical, default .F. |
| xlambda<X> | Interparticle-intra-level interaction <br> strength for non-spinors; <X>=1,2,3 | Real, default 0.d0. |


| xlambda12 | Interparticle-inter-level interaction <br> strength for non-spinors; | Real, default 0.d0 . |
| :--- | :--- | :--- |
| Spinor | Are the treated atoms spinors with <br> a spin-dependent interparticle inter- <br> action | Logical, default .F. |
| Lambda<X> | spin-independent (<X>=1) and spin- <br> dependent (<X>=2) interparticle in- <br> teraction strength, respectively, for <br> each level. | Real array, dimension 10, de- <br> fault 0.d0. |
| SpinOrbit | Toggles inclusion of spin-orbit- <br> interaction in the Hamiltonian | Logical, default .F. |
| Rashba_Prefactor | Magnitude of Rashba-spin-orbit- <br> interaction | Real, default 0.d0 |
| Dresselhaus_Prefactor | Magnitude of Dresselhaus-spin- <br> orbit-interaction | Real, default 0.d0 |
| SpinOrbit_Prefactor | Magnitude of total spin-orbit <br> (Rashba + Dresselhaus) term | Real, default 0.d0 |

It is important to note that it is necessary to specify at least as many one-body potentials as there are levels or spinor components in the treated particles. In the case of atoms featuring a conical intersection, the number of potentials is NLevel +1 . This is done in the routine Get_NLevelPotentials in the Get_1bodyPotential.F source file. The available predefined multilevel potentials are collected in the following table 4.2.

| whichpot | Description | Potential | Parameters |
| :--- | :--- | :--- | :--- |
| H01D | Parabolic potentials with <br> different frequencies and off- <br> set for different spin compo- | $V_{\uparrow}(x)$ <br> $V_{\downarrow}(x)=\frac{1}{2} p_{2}^{2}\left(x-p_{3}\right)^{2}+p_{1}^{2} x^{2} ;$ | $p_{1}, p_{2}$ are the frequencies of <br> the $\uparrow, \downarrow$ components/levels, <br> respectively. $p_{3}$ is the hor- |
| nents or levels. This poten- |  |  |  |
| inontal displacement of the |  |  |  |
| tial defined for two-level or |  |  |  |
| spin- $-\frac{1}{2}$ atoms, only. |  |  |  |$\quad$| two parabolas and $p_{4}$ their |
| :--- |
| relative offset. |

## Special parameters to treat ultracold atoms in an optical cavity

To deal with a system of bosons in interaction with an optical cavity, i.e., a field of photons which in turn generates a one-body potential for the atoms through the dipole force, special input parameters have been defined. These are listed in the following table 4.2.

| System Parameters Namelist |  |  |
| :---: | :---: | :---: |
| Parameter | Meaning | Options |
| Cavity_BEC | Toggle cavity treatment | Logical, default .F. |
| NCavity_Modes | How many modes to take into account treating the cavity | Integer, default 1 |
| Cavity_PumpRate | Rate, at which the laser is pumping the cavity (through the atoms) | Real, default 0.d0 |
| Cavity_LossRate | Rate, at which photons are lost from the cavity | Real, default 0.d0 |
| Cavity_K0 | Magnitude of the wave-vector defining the resonance frequency of the cavity | Real, default 0.d0 |
| Cavity_AtomCoupling | Coupling strength of the atomic resonance to the pumping laser frequency | Real, default 0.d0 |
| Pump_Switch | Is the pumping laser intensity ramped up exponentially, kept constant and then ramped down exponentially? | Logical, default .F. |
| RampupTime | Time over which the pump laser intensity is increased linearly up to Cavity_PumpRate. | Real, default 0.d0 |
| RampdownTime | Time over which the pump laser intensity is kept constant at Cavity_PumpRate. | Real, default 0.d0 |
| PlateauTime | Time over which the pump laser intensity is decreased linearly to 0.d0. | Real, default 0.d0 |
| Pump_Oscillate | Does the pump power oscillate [as $\left.\epsilon \sin \left(\omega_{p} t\right)\right]$ when it reached the plateau in the exponential ramping procedure? | Logical, default .F. |
| Pump_Amplitude | Amplitude $\epsilon$ of the oscillation of the pump power in units of Cavity PumpRate. | Real, default 0.d0 |
| Pump_Period | Period $\omega_{p}$ of the oscillation of the pump power. | Real, default 0.d0. |
| X_Cavity_Pump_Waist | Gaussian envelope's width for pump laser for two-dimensional systems | Real, default 0.d0 |
| Cavity_Mode_Waist | Gaussian envelope's width for cavity mode in two-dimensional systems | Real, default 0.d0 |
| Error_Rescale_Cavity | Parameter in the integration namelist to rescale the error tolerance of the integrator of the cavity equation of motion | Real, default 1.d0 |

## Special parameters to treat atoms in optical lattices

MCTDH-X offers two ways of dealing with atoms in optical lattices. Lattice Hamiltonians can be treated using an exact diagonalization treatment for one-, two-, and three-dimensional lattices. Especially, in the two- and three-dimensional cases, the dimensionality of the Hilbert space and the matrix that has to be diagonalized in the exact diagonalization approach is exploding rapidly and the problem size can no longer be handled. For the larger systems, one has to use the so-called MCTDBH approach, where the Bose-Hubbard Hamiltonian is expanded in a multiconfigurational basis with a number of effective single-particle states than there is lattice sites. With this approach, a systematic improvement beyond mean-field theories such as the discrete non-linear Schrödinger equation is possible. Contrary to the approaches like time-evolved block decimation, matrix product states or time-dependent density-matrix renormalization group, MCTDBH can provide accurate predictions in time and for higher dimension than $D=1$, because the partitioning of the Hilbert space is done by the variational principle and not artificially introduced. The parameters to select the former exact diagonalization or the latter MCTDBH approach are collected in the following table 4.2 .

| System Parameters Namelist |  |  |  |
| :--- | :--- | :--- | :---: |
| Parameter | Meaning | Options |  |
| Bose_Hubbard | Do an exact diagonalization of a <br> Bose Hubbard Hamiltonian? Atten- <br> tion: Has to be used together with <br> JOB_TYPE $=$ 'FCI' ! | Logical, default .F. |  |
| Periodic_BH | Has the treated lattice system peri- <br> odic boundary conditions? | Logical, default .F. |  |
| DVR Parameters Namelist |  |  |  |
| DVR_<I> | If DVR_<I> = 6, MCTDHB is ap- <br> plied to a lattice of that many sites <br> in I = X/Y/Z direction. Atten- <br> tion: has to be used together with <br> JOB_TYPE = 'BOS' ! |  |  |

It is important to note that in the case that an exact diagonalization of a Bose-Hubbard Hamiltonian is done, the one-body potential offset is obtained from the routine Get_BH_Offset. If MCTDHB applied to the Bose-Hubbard Hamiltonian the default routine Get_1bodyPotential is used to compute the on-site potential energy offset. Both routines can be found in the file ./source/ini_guess_pot/Get_1bodyPotential.F .

## The input of the analysis program

The analysis.inp file contains the desired quantities of analysis and specifies, for which points in time these are needed. and table 8 for all currently available parameters in the analysis program and their meaning.

| Parameter | Meaning | Options |
| :---: | :---: | :---: |
| ZERO_body |  |  |
| Time_From | Time from which to start analysis. | Real, no default. |
| Time_to | Time at which to stop analysis. | Real, no default. |
| Time_Points | Time points in analysis. | Integer, no default. |
| Total_Energy | Comupte the total, kinetic, potential and interaction energies | Logical, default .F. |
| Orbitals_Output | Create ASCII orbital output. | Logical, default .T. |
| FTOrbitals_Output | Create ASCII Fourier-transformed orbital output. | Logical, default .F. |
| Coefficients_Output | Create ASCII coefficients output. | Logical, default .T. |
| MatrixElements_Output | Output of reduced one-body and two-body density matrix elements. | Logical, default .F. |
| HamiltonianElements_Output | Output of one-body and two-body Hamiltonian matrix elements. | Logical, default .F. |
| GetA | Partial sums on the two-body Hamiltonians' matrix elements | Logical, default .F. |
| Dilation | Factor to dilate real space in order to obtain better k-space resolution (Only touched if AutoDilation is false) | Integer, default 1 |
| AutoDilation | Toggle if threshold Kdip is used to evaluate optimal dilation for FFTs (only for one-dimensional computations!) | Logical, default .F. |
| Kdip | Threshold to compute optimal dilation for FFTs if AutoDilation is true | Real, default 0.0001 |
| ONE_body |  |  |
| Density_x | Output of diagonal of one-body density in space. | Logical, default .F. . |
| Density_k | Output of diagonal of one-body density in momentum space. | Logical, default .F. |
| Pnot | Computation of nonescape probability $P_{\text {not }}$ ? | Logical, default .F. |
| xstart | Where does the integration on the density for $P_{\text {not }}$ start? | Real, no default. |
| xend | Where does the integration on the density for $P_{n o t}$ stop? | Real, no default. |
| ystart | Where does the integration on the density for $P_{\text {not }}$ start? | Real, no default. |
| yend | Where does the integration on the density for $P_{n o t}$ stop? | Real, no default. |
| zstart | Where does the integration on the density for $P_{\text {not }}$ start? | Real, no default. |


| zend | Where does the integration on the density for $P_{n o t}$ stop? | Real, no default. |
| :---: | :---: | :---: |
| Phase | Computation of the phase? | Logical, default .F. |
| Gradient | Computation of the phase gradient? | Logical, default .F. |
| Cavity_Order | Computation of cavity order parameter | Logical, default .F. |
| TWO_body |  |  |
| Correlations_X | Computation of spatial correlation functions on the full grid? | Logical, default .F. |
| Correlations_K | Computation of momentum correlation functions on the full grid? | Logical, default .F. |
| StructureFactor | Output of dynamic structure factor and local correlation functions | Logical, default .F. |
| xref | x value of reference point of dynamic structure factor | Real, default 0.d0 |
| yref | y value of reference point of dynamic structure factor | Real, default 0.d0 |
| zref | $z$ value of reference point of dynamic structure factor | Real, default 0.d0 |
| Correlation_Coefficient | Output of correlation coefficient $\tau=$ $\frac{\left\langle\vec{r}_{1} \vec{r}_{2}\right\rangle-\left\langle\left.\overrightarrow{r_{1}}\right\|^{2}\right.}{\left\langle\vec{r}^{2}\right\rangle-\langle\vec{r}\rangle^{2}}$ | Logical, default .F. |
| Geminals | Toggle output of natural geminal occupations in GO_PR.out | Logical, default .F. |
| FullGeminals | Toggle output of natural geminals in <time>NzMy-x-geminals.dat files | Logical, default .F. |
| corr1restr | Computation of spatial first order correlation functions on a restricted grid? | Logical, default .F. |
| xini1 | Restricted grid start | Real, no default. |
| xfin1 | Restricted grid stop | Real, no default. |
| xpts1 | Number of grid points for restricted grid. | Integer, no default. |
| corr1restrmom | Computation of spatial correlation functions on a restricted momentum grid? | Logical, default .F. |
| kxini1 | Restricted grid start | Real, no default. |
| kxfin1 | Restricted grid stop | Real, no default. |
| kpts1 | Number of grid points for restricted grid. | Integer, no default. |
| corr2restr | Computation of spatial second order correlation functions on a restricted grid? | Logical, default .F. |
| xini2 | Restricted grid start | Real, no default. |
| xfin2 | Restricted grid stop | Real, no default. |


| xpts2 | Number of grid points for restricted grid. | Integer, no default. |
| :---: | :---: | :---: |
| corr2restrmom | Computation of momentum second order correlation functions on a restricted grid? | Logical, default .F. |
| kxini2 | Restricted grid start | Real, no default. |
| kxfin2 | Restricted grid stop | Real, no default. |
| kpts2 | Number of grid points for restricted grid. | Integer, no default. |
| MANY_body |  |  |
| lossops | Loss operators, i.e., projectors on $N=2$ Hilbert space. | Logical, default .F. |
| border <br> Entropy | Border partitioning $N=2$ Hilbert space for evaluation of the loss operators. <br> Computation of diverse entropy measures | Real, no default. <br> Logical, default .F. |
| NBody_C_Entropy | Computation of entropy of matrix elements of full $N$-body density matrix | Logical, default .F. |
| TwoBody_Entropy | Computation of entropy of 2-body density matrix | Logical, default .F. |
| SingleShot_Analysis | Toggle the output of random deviates of the N-body density matrix. | Logical, default .F. |
| SingleShot_FTAnalysis | Toggle the output of random deviates of the N -body momentum density matrix. | Logical, default .F. |
| NShots | Number of single shots (random deviates of the N -body density) to compute. | Integer, default 10 |
| CentreOfMass | Toggle sampling of centre-of-mass operator | Logical, default .F. |
| CentreOfMomentum | Toggle sampling of centre-ofmomentum operator | Logical, default .F. |
| ShotVariance | Toggle computation of integrated variance of single shots | Logical, default .F. |
| NSamples | Number of samples to make from the centre-of-mass operator | Integer, default 10000 |
| anyordercorrelations_X | toggle output of higherorder <br> $\rho^{(p)}\left(\vec{r}_{\text {ref }}, \ldots, \vec{r}_{\text {ref }}, \vec{r}_{\text {order }-1}, \vec{r}_{\text {order }-2}\right)$ | Logical, default .F. |
| anyordercorrelations_X | toggle output of $\quad$ higher  <br> order correlations <br> $\rho^{(p)}\left(\vec{r}_{\text {ref }}, \ldots, \vec{r}_{\text {ref }}, \vec{r}_{\text {order }-1}, \vec{r}_{\text {order }-2}\right)$  | Logical, default .F. |


| order | specify order up to which correlation functions are to be output | Integer, default 10 |
| :---: | :---: | :---: |
| oneD | toggle output of correlations with one free variable $\vec{r}_{\text {order }}$ | Logical, default .T. |
| twoD | toggle output of correlations with two free variables $\vec{r}_{\text {order }-1}$ and $\vec{r}_{\text {order }}$ | Logical, default .F. |
| c_ref_x | x value of reference point for higher order correlations | Real, default 0.d0 |
| c_ref_y | y value of reference point for higher order correlations | Real, default 0.d0 |
| c_ref_z | z value of reference point for higher order correlations | Real, default 0.d0 |
| TWO_D |  |  |
| MOMSPACE2D | Output of 2D correlation functions' $g^{(1)}\left(\vec{r}_{1}^{\prime} \mid \vec{r}_{1}\right)$ and $g^{(2)}\left(\vec{r}_{1}^{\prime}, \vec{r}_{1}\right)$ in slices? | Logical, default .F. |
| REALSPACE2D | Output of 2D momentum correlation functions' slices? | Logical, default .F. |
| REALSKEW2D | Output of 2D skew correlation function $g^{(1)}(\vec{r} \mid-\vec{r} ; t)$ and $g^{(2)}\left(\vec{r}_{1},-\vec{r}_{1} ; t\right)$ ? | Logical, default .F. |
| MOMSKEW2D | Output of 2D skew momentum correlation function $g^{(1)}(\vec{k} \mid-\vec{k} ; t)$ and $g^{(2)}\left(\vec{k}_{1},-\vec{k}_{1} ; t\right)$ ? | Logical, default .F. |
| x1const | Keep X-coordinate of first position (momentum) $\left.\vec{r}_{1}\left(\vec{k}_{1}\right)\right)$ constant? | Logical, default .T. . |
| x1slice | At which value to keep the Xcoordinate of $\left.\vec{r}_{1}\left(\vec{k}_{1}\right)\right)$ ? | Real, default 0.0. |
| y1const | Keep Y-coordinate of first position (momentum) $\left.\vec{r}_{1}\left(\vec{k}_{1}\right)\right)$ constant? | Logical, default .T. . |
| y1slice | At which value to keep the Ycoordinate of $\left.\vec{r}_{1}\left(\vec{k}_{1}\right)\right)$ ? | Real, default 0.0. |
| x2const | Keep X-coordinate of second position (momentum) $\left.\vec{r}_{1}^{\prime}\left(\vec{k}_{1}^{\prime}\right)\right)$ constant? | Logical, default .F. |
| x2slice | At which value to keep the Xcoordinate of $\left.\vec{r}_{1}^{\prime}\left(\vec{k}_{1}^{\prime}\right)\right)$ ? | Real, default 0.0. |
| y2const | Keep Y-coordinate of second position (momentum) $\left.\vec{r}_{1}\left(\vec{k}_{1}\right)\right)$ constant? | Logical, default .F. |
| y2slice | At which value to keep the Ycoordinate of $\left.\vec{r}_{1}^{\prime}\left(\vec{k}_{1}^{\prime}\right)\right)$ ? | Real, default 0.0. |
| PROJ_X | Calculate effective density onedimensional potential $V_{\text {eff }}=$ $\int d \xi \sum_{i j} \rho_{i j} \phi_{i}^{*}(x, y, t) \phi_{j}(x, y, t)$, where $\xi=x$ and/or $y$ | Logical, default .F. |


| DIR | Specifying the direction for the effec- <br> tive one-dimensional potential. | Character, 'X' means $\xi=x$, <br> 'Y' means $\xi=y$ and 'B' <br> means both $\xi=x$ and $\xi=y$ <br> are computed. |
| :--- | :--- | :--- |
| L_Z | Computation of angular momentum <br> eigenvalue and matrix elements. | Logical, default .F. . |

Table 8: MCTDH-X analysis parameters

This documentation of the variables in the analysis.inp is also available in the html code documentation and the in-line documentation of the example file. After an appropriate modification, the program can be run (in an interactive shell) by typing

## MCTDHX

or by using/adapting/submitting one of the example PBS scripts in the PBS_Scripts directory and submitting the computation into the queue of a job scheduling system.

The analysis can be run after a computation has finished and the modification of the analysis.inp file with the following command:

```
MCTDHX_analysis
```

After the analysis program terminated successfully, ASCII files (structured as specified in the file documentation/Analysis_output_documentation) are in the working directory of the program. These files can be visualized using e.g. gnuplot or any other visualization software for data.

### 4.3 Available Visualization Scripts

The directory Visualization_Scripts contains several bash scripts to process the ASCII output of an MCTDHX_analysis run into .mpg or .avi movies. For a list of the currently available scripts and their function, see Table 9.

The way in which this is achieved, is by first processing the output files of an an MCTDHX_analysis run with mctdhx_gnuplot and produce a time-series of images. Subsequently, this time-series of images is encoded as a movie file using mctdhx_mencoder. mctdhx_gnuplot and mctdhx_mencoder are installed on your platform from the source tarball in the subdirectory External_Software by the MCTDH-X installation script. The automated way of using the visualization scripts is through the visualization master script. The Visualisation Master Script (vms.sh) is a short bash script that can be used to conveniently facilitate the creation of videos from a computation's data. To use vms.sh, just run it with the command

```
vms.sh <Movie Type> <Computation directory> <Lower plot range>
    <Upper plot range> <#Gridpoints> <2D-Slice 1> <2D-Slice 2>
```

where Movie type is one of the scripts in the Visualisation_Scripts directory. If Movie type=' 'all'' is specified, all available movie scripts will be run for the particular computation. The computation directory must contain the input file MCTDHX.inp of the computation. The other arguments are optional and need not to be given. They can be used to fine-tune the output of the

| Visualization Script | Function |
| :--- | :--- |
| 1D-CORR1-K | Movie of the coherence $\left\|g^{(1)}\right\|^{2}$ |
| 1D-CORR2-RESTR-X | Movie of the two-body correlations $\left\|g^{(2)}\right\|$ on a restricted <br> grid. |
| 1D-DENSITY_X-TIME-EVOLUTION-PM3D | "Movie" of the density $\rho(x, t)$, i.e., 2D plot where the <br> $y$-axis is time. |
| 1D-CORR1-RESTR-K | Movie of the coherence $\left\|g^{(1)}\right\|^{2}$ in momentum space on a <br> restricted grid. |
| 1D-CORR2-X | Movie of the two-body correlations $\mid g^{(2)}$. |
| 1D-CORR1-RESTR-X | Movie of the coherence $\left\|g^{(1)}\right\|^{2}$ on a restricted grid. |
| 1D-DENSITY_K | Movie of the momentum density $\rho(k, t)$. |
| 1D-CORR1-X | Movie of the coherence $\left\|g^{(1)}\right\|^{2}$. |
| 1D-DENSITY_X | Movie of the density $\rho(x, t)$. <br> 2D-DENSITY_M2 <br> dimensional computations. |
| 2D-DENSITY_M3 | Movie of the density and the first three orbitals in two- <br> dimensional computations. |
| 2D-DENSITY_M4 | Movie of the density and the first four orbitals in two- <br> dimensional computations. |
| 2D-ORB_AVG_PHASE_DENSITY_Lz | Movie of Orbital average phase, density and angular mo- <br> mentum in two-dimensional computations. |
| 2D-AVG-PHASE | Movie of the average phase in two-dimensional compu- <br> tations. |
| 2D-DENSITY_K | Movie of the momentum density in two-dimensional <br> computations. |
| 2D-AVG-PHASE-PHASEGRADIENT- <br> DENSITY-ENERGY-LZ | Movie of the average phase, phase gradient, density and <br> angular momentum in two-dimensional computations. |
| 2D-DENSITY_X | Movie of the density in two-dimensional computations. |
| CI | Movie of the CI coefficients. |
| nat_occ_loop | Plotting the natural occupations of a computation. |

Table 9: List of available visualization scripts.
All the above scripts take 5 command line arguments: start time, stop time, time increment, number of orbitals, number of particles. To generate a movie of the coherence on a restricted grid in momentum space one could do 1D-CORR1-RESTR-K 01000.14101 - this command would generate a movie for the first 100 time units in steps of 0.1 for an $M=4$ computation with $N=101$ particles. The visualization master script, vms.sh, automates the usage of the above movie scripts.
analysis program, i.e., the plotting range, the gridpoints and the slices for the output of the correlation functions of two-dimensional computations. If you need to find out the valid strings to use, then run the command $. / \mathrm{vms} . \operatorname{sh}$ and list of the available strings will be output to the screen. If the command ./vms.sh <Movie Type> <Computation directory> ... is entered successfully, it will generate all the ASCII data from the binary MCTDH-X output and build the movie(s). All the visualization scripts in the Visualization_Scripts subdirectory have five command line arguments: timeInitial, the time at which the movie will begin; timeIncrement, the time difference between each frame; timeFinal, the time at which the movie will end; MOrbs, the total number of orbitals in the simulation and Npar, the number of particles in the simulation. vms.sh basically calls the respective desired visualization script(s) with the appropriate arguments. Eventually, it prints a short message 'And we are happy.' once the video has been successfully created. Of course, one can also call the scripts in the Visualization_Scripts directory manually by typing

```
<Visualization_Script> <timeInitial> <timeFinal> <timeIncrement> <MOrbs> <Npar>
```

where <Visualization_Script> can be chosen from the list in Table 9.

### 4.4 Configuring the Monster Script

This script basically acts as a secretary, creating subdirectories, copying files, editing text, and submitting jobs in a coordinated fashion so that after the user configures just 2 text files, a single function call can result in thousands of computers working for days to weeks to calculate up to a combined 10-dimensional paramter scan. There are two steps to configuring MonsterScript.sh: configuring the input file and preparing the run files. An example input file can be found at Input_Examples/ParameterScan.inp which contains in-line documentation (see 10). Inside the working directory, one must place a properly configured ParameterScan.inp, libmctdhx.so, a binary executable file consistent with \$binary defined inside ParameterScan.inp, and an input template consistent with \$Input_Template defined inside ParameterScan.inp. Once these 4 files are properly configured and placed, the script is called by running
\$MCTDHXDIR/MonsterScript.sh. The script scans up to 5 user defined relaxation parameters, runs them until convergence is detected and, if desired, automatically scans each relaxation with up to 5 user-defined propagation parameters. When running on a cluster, the script will automatically restart jobs if they finish, due to time constraints, before the calculation is complete. To circumvent job number restrictions on clusters, the script will build a series of runscripts that each simultaneously run many calculations in a single job, rather than a single computation per job.

A little under-the-hood knowledge is useful to to effectively use and debug MonsterScript.sh. The script initially calls \$MCTDHXDIR/Computation_Scripts/ParameterScan_Propagation.sh or \$MCTDHXDIR/Comptuation_Scripts/ParameterScan_Relaxation.sh depending on the input file configuration, which then iterates through the corresponding parameter set. As the script loops through each parameter set, it calls either
\$MCTDHXDIR/Scripts/IterateParameters.sh or \$MCTDHXDIR/Scripts/IterateParameters_relax.sh which then perform the actual secretary functionality using about a dozen smaller scripts inside \$MCTDHXDIR/Scripts. If it is determined that a new calculation must be run when using a cluster, a few lines are added to a runscript in the working directory called run\#.sh for some number \#. To avoid duplicate computations, a file in the computation directory is created called RunFlag, which is automatically deleted when the job runs out of time or the computation is complete, via some code within the runscript. The runscript accumulates calculation jobs until the total number of nodes
requested reaches a threshold defined by \$MPMDNodes in ParameterScan.inp, and then it is submitted using (in most cases) a qsub command within either \$MCTDHXDIR/Scripts/MPMDrun_relax.sh or \$MCTDHXDIR/Scripts/MPMDrun_prop.sh on line 106. It is often useful to comment this line for testing purposes, as it will prevent the script from submitting any jobs, leaving it to only copy, move, and edit files. Directories with incomplete calculations are stored in files Relax_Array and Prop_Array*. Directories are removed from these files when their corresponding calculations are complete, triggering an end of the corresponding ParameterScan_*.sh function call when the array is empty.

When MonsterScript.sh is run without any arguments, all lingering runscripts, RunFlag's, and arrays are cleared. If the user wishes to keep these files, they must run MonsterScript.sh save.

| Parameter Name | Description | Values |
| :--- | :--- | :--- |
| Do_Relax | Specifies whether to do or skip relax- <br> ations. | "T", "F" |
| Do_Propagations | Specifies whether to do or skip propa- <br> gations | "T", "F" |
| Propagation_Start | Specifies how to start propagations <br> when relaxation is skipped | "BINR", "HAND" |
| Do_Analysis | Specifies whether to do or skip analysis | "T", "F" |
| MonsterName | Suffix for job names "MON- <br> STER_\$MonsterName" | any string |
| runhost | Specifies which cluster is used, or if no <br> cluster is used | "hermit", <br> "maia", "bwgrid", "PC" |
| numnodes | Specifies number of nodes used for re- <br> laxation jobs | Integer |
| MPMD | Specifies whether to run in Multiple <br> Program Multiple Data mode, i.e. mul- <br> tiple computations per submitted job. <br> This only applies to scans on a cluster, <br> and is recommended if more than 20 <br> computations are desired |  |
| MPMDjobs | If using MPMD mode, this specifies <br> how many nodes are requested for each <br> job | Integer |
| maxjobs | Maximum number of jobs allowed on <br> the queue (20 for hornet and hermit | Integer |
| binary | Name of MCTDHX executable in work- <br> ing directory | usually "MCTDHX_intel" |
| Relaxation_Template | Name of input template for relaxations <br> in working directory | usually "MCTDHX.inp" |


| Propagation_Template | Name of input template for propagations in working directory (if identical to Relaxation_Template, the parameters adjusted for the relaxation will be copied, if not then the paramters adjusted for the relaxation will not be copied). | usually "MCTDHX.inp" |
| :---: | :---: | :---: |
| Relaxtime | Time to run relaxations for | Positive number |
| NParameters | Number of parameters to scan for relaxations | [1-5] |
| Parameter\# | \#=1-5, name of relaxation parameter \# | Any parameter found in MCTDHX.inp |
| List\# | $\#=1-5$, specifies if a list of values is used for relaxation parameter \# | "T", "F" |
| Parameter\#List | $\#=1-5$, if \$List\#=" T ", specifies parameter \# values to scan over | Array |
| Scan\#_Start | \#=1-5, if if \$List\#="F", speficies beginning of range of relaxation parameter \# to be scanned | number |
| Scan\#_Stop | $\#=1-5$, if if \$List\#="F", speficies end of range of relaxation parameter $\#$ to be scanned | number |
| Scan\#_Step | \#=1-5, if if \$List\#="F", speficies step of range of relaxation parameter \# to be scanned | number |
| MaxNodes | Maximum number of nodes allocated for a single propagation computation | Positive integer |
| Prop_Time_Final | End time of propagation computations | Positive number |
| Prop_NParameters | Number of propagation parameters to scan over | [1-5] |
| Prop_Parameter\# | \#=1-5, name of propagation parameter \#. | Any parameter in "MCTDHX.inp" |
| Prop_List\# | $\#=1-5$, specifies if a list of values is used for propagation paramter \# | "T","F" |
| Prop_Parameter\#List | \#=1-5, if \$Prop_List\#="T", specifies propagation parameter $\#$ values to scan over | list of appropriate values |
| Prop_Scan\#_Start | \#=1-5, if \$Prop_List\#=" $F$ ", specifies beginning of scan range for propagation parameter \# | Number |
| Prop_Scan\#_Stop | \#=1-5, if \$Prop_List\#="F", specifies end of scan range for propagation parameter \# | Number |
| Prop_Scan\#_Step | \#=1-5, if \$Prop_List\#="F", specifies step of scan range for propagation parameter \# | Number |

Table 10: Parameters inside ParameterScan.inp.

## 5 MCTDH-X main and analysis program output documentation

### 5.1 Main program output

MCTDH-X has several standard output files, and it will generate additional ASCII output if the toggle Write_ASCII is set to true in the input of a computation. Specifically, standard output is comprised by the files NO_PR.out, Initialization.dat, Error.dat, and Timing.dat. The additional ASCII output is comprised of <time>orbs.dat and <time>coefs.dat files - these can also be generated after a computation has finished by running the analysis program which processes the binary data files PSI_bin and CIc_bin. Set aside the binary PSI_bin and CIc_bin and the ASCII Initialization.dat files, all the above output files are column-formatted ASCII files.

### 5.1.1 Output structure in standard MCTDHX computations

For the case that atoms without internal structure are treated (Multi_Level = .F.), the column structure of the ASCII output files is described in the following tables 11,12,13, 14, 15 .

| Column 1 | Column 2 to $(1+M)$ | Column $2+M$ |
| :--- | :--- | :--- |
| Time $t$ | Natural orbital oc- <br> cupations $\rho_{M}^{(N O)}(t)$ to <br> $\rho_{1}^{(N O)}(t)$ | Energy $E(t)$ |

Table 11: NO_PR. out file structure.
This table explains the column structure of the NO_PR. out file. $M$ stands for the number of orbitals in the computation.

### 5.1.2 Output structure in block Davidson relaxations

In the case of block Davidson computations, several vectors of coefficients are relaxed simultaneously using the same set of orbitals. Consequently, the structure of the output changes as specified in the following tables 16, and 17 .

For every vector in the block, a separate file with natural occupations is generated whose structure is identical to table 11. The naming convention for these files is NO_PR.BL<state> where <state> stands for the index of the vector in the block.

### 5.1.3 Output structure in multilevel MCTDHX computations

For the case that atoms with internal structure are treated (Multi_Level = .T.), the column structure of the ASCII output files is described in the following tables $18 / 19$.

| Column 1 | Column 2 | Column 3 | Column 4 | Column 5 | Column 6 | Column 7 | Column 8 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathrm{N}^{o}$ of in- <br> tegration <br> step | Time $t$ | $E_{\text {orb }}(t)$ | $E_{\text {orb,rel }}(t)$ | $E_{C I}(t)$ | $E_{\text {CI,rel }}(t)$ | $E_{\text {tot }}(t)$ | $E_{V_{t}}(t)$ |

Table 12: Structure of the Error. dat file.
This table explains what is saved in the different columns of the Error.dat file. $E_{\text {orb }}(t)$ and $E_{\text {orb,rel }}(t)$ are the absolute and relative integration errors from the orbitals' equations of motion, respectively. $E_{C I}(t)$ and $E_{C I, r e l}(t)$ are the integration errors from the coefficients' equations of motion, respectively. $E_{\text {tot }}(t)=E_{\text {orb }}(t)+E_{C I}(t)$ is the sum of the orbital and coefficients integration errors and $E_{V_{t}}(t)$ it an (experimental) error measure for the error due to a time-dependency of the external one-body potential.

| Column 1 | Column 2 | Column 3 | Column 4 | Column 5 | Column 6 | Column 7 | Column 8 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| N of in- <br> tegration <br> step | Time $t$ | Execution <br> time | $T_{\text {step }}$ | $T_{C I}$ | $T_{\rho}$ | $T_{\text {CI,Func }}$ | $T_{\text {Orb }}$ |

Table 13: Timing. dat file structure.
this table displays the column structure of the Timing. dat output file. Here, $T_{\text {step }}$ is the execution time for the present integration step, $T_{C I}=T_{\rho}+T_{C I, F u n c}$ is the overall execution time in this step spent on the configuration interaction, i.e., the coefficients part of the program. $T_{\rho}$ is the runtime consumed to invert the matrix elements of the reduced one-body density, $T_{C I, F u n c}$ is the execution time consumed in applying the Hamiltonian to the coeffcients vector. $T_{\text {Orb }}$, finally, is the execution time spent for evaluating the right hand side of the orbitals' equations.

| Column 1 to 3 | Column 4 | Column 5 | Column 6 \& 7 | Column 8 \& 9 | Column 10 \& 11 to <br> $(10+2 M) \&(11+2 M)$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $x, y, z$ | DVR <br> weight | $V(x, y, z, t)$ | $\rho_{w}(x, y, z ; t)$ | $\rho_{(N O)}(x, y, z ; t)$ | $\phi_{1}(x, y, z ; t)$ to <br> $\phi_{M}(x, y, z ; t)$ |


| $\left.\left\lvert\, \begin{array}{l}\text { Column }(11+2 M+1) \\ \&(11+2 M+2) \\ \text { to } \\ (11+4 M+1)\end{array}\right.\right)(11+$ |
| :--- |
| $4 M+2)$ |
| $\phi_{M}^{(N O)}(x, y, z ; t) \quad$ to |
| $\phi_{1}^{(N O)}(x, y, z ; t)$ |

Table 14: <time>orbs. dat file structure.
This table explains the column structure of the <time>orbs. dat output files of the main or analysis program. $x, y, z$ are the spatial coordinates, $V(x, y, z, t)$ is the one-body potential, $\rho_{w}(x, y, z ; t)$ is the density in working orbitals, $\rho_{(N O)}(x, y, z ; t)$ is the density in natural orbitals, $\phi_{1}(x, y, z ; t)$ to $\phi_{M}(x, y, z ; t)$ are the working orbitals, and $\phi_{M}^{(N O)}(x, y, z ; t)$ to $\phi_{1}^{(N O)}(x, y, z ; t)$ are the natural orbitals. Please note, that some of the quantities are complex numbers which then are output decomposed in their real and imaginary parts in two columns (as specified by the column numbers).

| Column 1 | Column 2 \& Column 3 |
| :--- | :--- |
| $\mathrm{N}^{o}$ of Coefficient | Real and imaginary part of <br> the coefficient |

Table 15: Structure of the <time>coef.dat files

| Column 1 | Column 2 to Column 2 • blocksize +1 |
| :--- | :--- |
| $\mathrm{N}^{o}$ of Coef- <br> ficient | Real and imaginary part of the coeffi- <br> cient for each vector in the block |

Table 16: Structure of the <time>coef. dat files for block Davidson computations

| Column 1 to 3 | Column 4 | Column 5 | Column 6 \& 7 | Column 8\& 9 <br> to $(8+2 M) \&$ <br> $(9+2 M)$ | $\cdots$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $x, y, z$ | DVR <br> weight | $V(x, y, z, t)$ | $\rho_{(N O)}(x, y, z ; t)$ | $\phi_{M}^{(N O)}(x, y, z ; t)$ <br> to <br> $\phi_{1}^{(N O)}(x, y, z ; t)$ | Repeat columns <br> $6 \cdots 9+2 M$ for all <br> vectors in the block |

Table 17: <time>orbs.dat file structure for block Davidson computations.
This table explains the column structure of the <time>orbs. dat output files of the main or analysis program in the case of Block-Davidson relaxations. $x, y, z$ are the spatial coordinates, $V(x, y, z, t)$ is the one-body potential, $\rho_{(N O)}(x, y, z ; t)$ is the density in natural orbitals, and $\phi_{M}^{(N O)}(x, y, z ; t)$ to $\phi_{1}^{(N O)}(x, y, z ; t)$ are the natural orbitals. Please note, that some of the quantities are complex numbers which then are output decomposed in their real and imaginary parts in two columns (as specified by the column numbers). Importantly, the dots $\cdots$ imply that the columns 6 to $9+2 M$ are repeated for all the wavefunctions in the block Davidson computation.

| Column 1 | Column 2 to $(1+M)$ | Column $(2+M)$ | Column $(3+M)$ to $3+M+N_{l}$ |
| :--- | :--- | :--- | :--- |
| Time $t$ | Natural orbital occupations <br> $\rho_{M}^{(N O)}(t)$ to $\rho_{1}^{(N O)}(t)$ | Energy $E(t)$ | State populations (density) for <br> all levels |

Table 18: NO_PR.out file structure for computations with Multi_Level=.T..
This table explains the column structure of the NO_PR. out file. $M$ stands for the number of orbitals in the computation and $N_{l}$ is the number of internal states considered.

| Column 1 to 3 | Column 4 | Column 5 to | Column 5 + | Column 5 + |
| :--- | :--- | :--- | :--- | :--- |
|  |  | $4+N_{l}+N_{C I}$ | $N_{l}+N_{C I} \&$ | $3 N_{l}+N_{C I} \&$ |
|  |  |  | $6+N_{l}+N_{C I}$ to | $6+3 N_{l}+N_{C I}$ |
|  |  |  | $3+3 N_{l}+N_{C I}$ | to $3+5 N_{l}+$ |
|  |  |  | $\& 4+3 N_{l}+$ | $N_{C I} \& 4+$ |
|  |  |  | $N_{C I}$ | $5 N_{l}+N_{C I}$ |
| $x, y, z$ | DVR weight | $V^{i}(x, y, z, t)$ | $\rho_{w}^{i}(x, y, z ; t)$ | $\rho_{(N O)}^{i}(x, y, z ; t)$ |
|  |  | and |  |  |


| Column $5+5 N_{l}+N_{C I} \& 6+5 N_{l}+$ | Column $5+5 N_{l}+N_{C I}+2 M N_{l} \&$ |
| :--- | :--- |
| $N_{C I}$ to $3+5 N_{l}+N_{C I}+2 M N_{l} \&$ | $6+5 N_{l}+N_{C I}+2 M N_{l}$ to $3+5 N_{l}+$ |
| $4+5 N_{l}+N_{C I}+2 M N_{l}$ | $N_{C I}+4 M N_{l} \& 4+5 N_{l}+N_{C I}+$ |
|  | $4 M N_{l}$ |
| $\phi_{k}^{i}(x, y, z ; t)$ for $i=1, . ., N_{l}$ and | $\phi_{k}^{(N O), i}(x, y, z ; t)$ for $i=1, . ., N_{l}$ |
| $k=1, . ., M$ | and $k=M, . .1$ |

Table 19: <time>orbs.dat file structure for multilevel computations.
This table explains the column structure of the <time>orbs. dat output files of the main or analysis program for the case that Multi_Level=.T. was set. $N_{l}$ is the number of internal states and $N_{C I}$ is the number of conical intersections ( $N_{C I}=0$ if Conical_Intersection=.F.). Please note, that the index of the internal state is always running first before the orbitals' index. $x, y, z$ are the spatial coordinates, $V^{i}(x, y, z, t)$ is the one-body potential of internal state $i, V_{C I}^{i}(x, y, z, t)$ is the coupling of the $i$-th conical interaction, $\rho_{w}^{i}(x, y, z ; t)$ is the density in working orbitals for internal stat $i, \rho_{(N O)}^{i}(x, y, z ; t)$ is the density in natural orbitals for state $i, \phi_{1}^{i}(x, y, z ; t)$ to $\phi_{1}^{M}(x, y, z ; t)$ are the working orbitals in internal state $i$, and $\phi_{M}^{(N O), i}(x, y, z ; t)$ to $\phi_{1}^{(N O), i}(x, y, z ; t)$ are the natural orbitals in state $i$. Please note, that some of the quantities are complex numbers which then are output decomposed in their real and imaginary parts in two columns (as specified by the column numbers).

The structure of the output files described in tables 12 13 15, is identical for multilevel computations.

### 5.2 Analysis program output

The output of the MCTDH-X analysis software is toggled with the input file analysis.inp as described in the table 8. Generally there are two kinds of different output file structures. Some analysis quantities are scalar and their time series will be saved in a single file, like e.g., the nonescape probability and some other quantities need one file per point in time, like e.g., the density or momentum density. Orbitals_Output and Coefficients_Output toggle the output of
the files <time>orbs.dat and <time>coefs.dat - their structure is the same as when one sets the variable Write_ASCII in the main programs input to .T., see tables 14 and 15 .

The structure of the output files generated by setting the respective analysis variables is collected in Appendix ?? in tables 22|23|24|25|26|27|28||29|30|31|32|33|34|35|36, and 37 .

## 6 Developer Guidelines

The MCTDH-X program package uses Mercurial as a version management and provides a documentation in .pdf format as well as a .html code documentation generated by Doxygen. Generally, if you are planning to implement something useful it is recommended to browse the code documentation in documentation/html/index.html to find out in which module to start. The package developers will also be happy to help you with this.

In writing code, please stick to the following principles:

1. develop code in different branches of the repository (see section 7 below for how to create/manage branches)
2. write code that is readable (use indentation and obvious variable and subroutine names)
3. write code that is modular (group written subroutines or variable declarations into modules)
4. use CamelCase naming conventions, where appropriate. Every new word in declarations (of subroutines, variables or modules) should start with a capital letter. For subroutines and modules their action(s) are prefixed with an underscore, like for instance Get_KineticEnergyAction_AllOrbitals.
5. update this manual with new functionality and user guidance
6. add code documentation that doxygen can process, i.e., start commented lines with !> or !<. This is especially crucial for new subroutines, such that other users are able to understand how your routine works and may use or further develop it.
7. share your version by commiting it to your working copy and letting the developers know, such that your development branch can be merged and tested to enter the next release of the package (see paragraph on the version management below).
8. use the test-script testMCTDHX.sh to test your developments against reference values.

## 7 Version Management

The tool that is used to manage the development of the MCTDH-X software is Mercurial, in terminal, hg. It is available on most Unix-based systems and facilitates the contribution of multiple developers to a project. Mercurial is a distributed version management system, i.e., each user of it has the the full version history available locally. The basic operations are collected in the following list:

1. making a full copy of a repository with hg clone <repository location>
2. creating branches with hg branch <branch name>.
3. showing a summary of the changes with respect to the last revision by hg status
4. adding and removing files by hg add <files> and hg remove <files>
5. add all new files and remove all missing files with hg addremove
6. showing the history of commit messages for all revisions by hg log
7. add a version number to your revision with hg tag <version label>. The structure of the version label is explained below.
8. commiting changes made to a revision of the repository as a new revision with hg commit -u <username>
9. pushing or pulling the changes made in the present repository to another repository by hg push <repository location> or hg pull <repository location>.
10. updating the repository after pulling/pushing a changeset by hg update

It is important to note here that the version management of the central repository is not open for all users. If you are interested in contributing, write an email to the developers at mctdhx@ultracold.org to obtain an account for the software repository.

Version labels: The version labels for MCTDH-X follow the following syntax:

```
<major>.<minor><small><fix>
```

Here, <major> is an integer that is incremented by 1 when a new major feature is commited (for instance the capability of the program to treat mixtures of indistinguishable particles). The <minor> label is an integer, that is incremented by 1 when a new minor feature like, for instance, a new operator or a new analysis routine is ready. The small> label is an integer which is incremented by 1 , whenever changes to the repository are made which do not correspond to a feature like, for instance, an improvement of the efficiency or an update to the manual.

## Contribution how-to:

The strategy taken for the development of MCTDH-X is that every contributor should create her/his own branch(es) for the software development. Branches may be created with a name describing the respective contribution or just the contributor's name. When a new version is due to be released, all the different branches will be merged into the default branch to form a release.

In the following, the basic steps to create branches and to push your current revision into the main repository are described. These include to first clean up your current version, second make sure if the changes you want to share as a contribution to the main repository are there and which files are affected, third commit the changes as a new revision to your local repository and write a thorough description of the implemented changes and fourth, push the changeset to the main repository.

Create branch: To create a branch, run hg branch <name>. Typically <name> should be your user name in the forum or descriptive for the features that are planned to be developed in this branch.

Clean up: To remove the compiled included software, run make -f <Makefile.your-configuration> purge and then, to remove the objects and libraries generated in the compilation of the main program, run make -f <Makefile.your-configuration> clean.

Inspect the changeset: Check if indeed only the files that you intended to modify show up when you run hg status. If unsure, look through the individual files and in case of doubt, recompile/reinstall the program and test it again (for instance by running testMCTDHX.sh) and start over with the first step.

Commit revision: Once you verified and tested your changes, commit them to your local repository by hg commit -u <username>. After issueing this command you will be prompted to write a so-called commit message. The text should include a brief but concise description of the changes entering the revision and a list of open tasks or known problems. The changelog of the program is the timeline of these commit messages and they are therefore crucial to let the other users and programmers know about changes.

Share the changes: After your local revision is prepared, you can share the changeset by contributing it to the main repository by running hg push, which will prompt you for username and password. If you get a message that notifies you of the creation of a remote head, please contact the developers at mctdhx@ultracold.org since it's likely that there is a conflict between your changeset and the latest revision in the repository. In that case, it is necessary to merge the heads and resolve the conflicts.

## Feedback and suggestions as well as bug-reports are welcome anytime at metdhx@ultracold.org or http://ultracold.org/forum.

## A Predefined one-body potentials

| whichpot | Description | Potential | Parameters |
| :---: | :---: | :---: | :---: |
| H01D | 1D harmonic oscillator | $V(x)=\frac{1}{2} p_{1}^{2} x^{2}$ | $p_{1} \equiv$ trap frequency. |
| H02D | 2D harmonic oscillator | $\begin{aligned} & V(x, y) \\ & \frac{1}{2}\left(p_{1}^{2} x^{2}+p_{2}^{2} y^{2}\right) \end{aligned}$ | $\begin{aligned} & p_{1 / 2} \equiv \operatorname{trap} \text { frequency } \\ & \text { in } x / y \end{aligned}$ |
| H03D | 3D harmonic oscillator | $\begin{aligned} & V(x, y, z) \\ & \frac{1}{2}\left(p_{1}^{2} x^{2}+p_{2}^{2} y^{2}+p_{3}^{2} x^{2}\right) \end{aligned}$ | $\begin{aligned} & \hline p_{1 / 2 / 3} \equiv \text { trap frequency } \\ & \text { in } x / y / z \\ & \hline \end{aligned}$ |
| h+d | 1D harmonic oscillator plus Gaussian central barrier | $\begin{aligned} & V(x)=\frac{1}{2}\left(x-p_{1}\right)^{2} \\ & +p_{2} \exp \left(-\frac{\left(x-p_{1}\right)^{2}}{2 p_{3}^{2}}\right) \end{aligned}$ | $\begin{aligned} & p_{1} \equiv \text { diplacement }, \\ & p_{2} \equiv \text { height of Gaus- } \\ & \text { sian, } p_{3} \equiv \text { width of } \\ & \text { Gaussian } \end{aligned}$ |
| h2D+d | 2D harmonic oscillator plus Gaussian central barrier | $\begin{aligned} & V(x)=\frac{1}{2}\left(p_{1}^{2} x^{2}+p_{2}^{2} y^{2}\right) \\ & +p_{3} \exp \left(-\left(\frac{(x)^{2}}{2 p_{4}^{2}}+\frac{(y)^{2}}{2 p_{5}^{2}}\right)\right) \end{aligned}$ | $p_{1 / 2} \equiv$ trap frequency in $x$ - $/ y$-direction, $p_{3} \equiv$ height of Gaussian, $p_{4 / 5} \equiv$ width of Gaussian in direction $x / y$ |
| HO<X>D+td_gauss | $<X>D$ harmonic oscillator plus time-dependent [height $A(t)$ ] Gaussian central barrier in $x$-direction | $\begin{aligned} & V(x)=\frac{1}{2} p_{1}^{2}\left(\left(x-p_{2}\right)^{2}+y^{2}\right) \\ & +A(t) \exp \left(-\frac{\left(x-p_{5}\right)^{2}}{2 p_{6}^{2}}\right) \end{aligned}$ | $p_{1} \equiv \operatorname{trap}$ frequency in $x$ - and $y$-direction, $p_{2} \equiv$ displacement of the minimum of the trap w.r.t. the barrier, $p_{5} \equiv$ displacement of Gaussian w.r.t. $x=0$, $p_{6} \equiv$ width of Gaussian barrier in direction $x$. Height $A(t)=$ $\begin{cases}\frac{t p_{3}}{p_{4}} & t \leq p_{4} \\ p_{3} & t>p_{4}\end{cases}$ |
| Tilt | Tilted triple well potential | $\begin{aligned} & V(x)=-p_{1} x+p_{2} \sin (2 x)^{4}+ \\ & (x / 2.2)^{20} \end{aligned}$ | $\begin{aligned} & p_{1} \equiv \text { tilt parameter, } \\ & p_{2} \equiv \text { depth of the } \\ & \text { wells. } \end{aligned}$ |
| Tiltinit | single well potential to initialize system for a propagation in the tilted triple well potential Tilt. | $\begin{aligned} & V(x)=V(x)=-p_{1} x+ \\ & p_{2} \sin (2 x)^{4}+\left[\left(x-\frac{p_{3} \pi}{2}\right) / 0.7\right]^{20} \end{aligned}$ | $p_{1} \equiv$ tilt parameter, $p_{2} \equiv$ depth of the wells, $p_{3}$ is used to select displacement from the origin. |
| OL_HW_1D | lattice in one dimension with hard wall boundaries | $\begin{aligned} & V(x) \\ & \begin{cases}p_{1} \sin \left(p_{2} x\right)^{2} & p_{3}<x<p_{4} \\ 1000 & \text { else }\end{cases} \end{aligned}$ | $p_{1} \equiv$ depth of lattice, $p_{2} \equiv$ frequency of lattice, $p_{3}, p_{4}$ are the boundaries for the hard walls. |


| TDHIM | Harmonic potential with time-dependent frequency for benchmarks with the time-dependent harmonic interaction Hamiltonian. | $\begin{aligned} & V(x)=\frac{1}{2}(1+\sin (t) \cos (2 t) \\ & \left.\sin \left(\frac{1}{2} t\right) \sin (0.4 t)\right) x^{2} \end{aligned}$ | No $\quad$ parame- ters, use with which_interaction ='TDHIM' |
| :---: | :---: | :---: | :---: |
| tun | Potential for tunneling to open space dynamics | $\begin{aligned} & V(x) \\ & \begin{cases}\frac{1}{2} x^{2} & x \leq 2 \\ 2.2662969 & \\ \exp \left(-2(x-2.25)^{2}\right) & x>2\end{cases} \end{aligned}$ | no parameters, initial wavefunction should be localized at $x=0$, use whichpot='ini' for relaxation of initial state. |
| thr | Tunneling to open space with a threshold | $\begin{aligned} & V(x) \\ & \begin{cases}\frac{1}{2} x^{2} & x \leq 2 \\ A x^{3}+B x^{2} & = \\ +C x+D & 2<x \leq 4 \\ p_{1} & x>4\end{cases} \end{aligned}$ | Parameter $p_{1}$ defines the threshold and the polynomial coefficients $A=1-\frac{1}{4} p_{1}$, $B=2.25 p_{1}-9.5$, $C=28-6 p_{1}$, $D=5 T-24$. |
| DQD | Double quantum dot potential | $\begin{aligned} & V(x)=-p_{1} \exp \left(-p_{2}(x+\right. \\ & \left.\left.\frac{1}{2} p_{3}\right)^{2}\right)-p_{4} \exp \left(-p_{5}\left(x-\frac{1}{2} p_{3}\right)\right) \end{aligned}$ | $p_{1} \equiv$ depth of first quantum dot, $p_{2} \equiv$ width of first quantum dot, $p_{3} \equiv$ distance of the two dots, $p_{4} \equiv$ depth of second quantum dot, $p_{5} \equiv$ width of second quantum dot. |
| DQDLASER | Double quantum dot potential illuminated by a laser with time-dependent amplitude $V_{\text {las }}(x, t)$. | $\begin{aligned} & V(x)=-p_{1} \exp \left(-p_{2}(x+\right. \\ & \left.\left.\frac{1}{2} p_{3}\right)^{2}\right)-p_{4} \exp \left(-p_{5}(x-\right. \\ & \left.\left.\frac{1}{2} p_{3}\right)\right)+V_{\text {las }}(x, t) \exp \left(-p_{9}(x-\right. \\ & \left.\left.p_{10}\right)^{2}\right) \end{aligned}$ | $p_{1-5} \equiv$ same as in DQD, $\quad V_{\text {las }}(x, t)=$ $p_{7} \cos \left(p_{8} t\right) \sin \left(\pi \frac{t}{p_{6}}\right) x$ for $t \leq p_{6}$ and $V_{\text {las }}(x, t)=0$ else. |
| zero | No potential | $V=0$ | No parameters. Boundary conditions determined by the discrete variable representation. |
| MQT_ini | Rectangular 1D box | $V(x)=\infty \quad \forall[x \quad \notin$ $\left.\left(0,20-p_{1}\right)\right]$ and $V(x)=$ $0 \quad \forall[x \in) 0,20-p_{1}(]$ | $p_{1} \equiv$ barrier width in propagation MQT_prop |
| MQT_prop | Rectangular 1D box, barrier and open space | $\begin{aligned} & V(x)=\infty \quad \forall[x<0] \text { and } \\ & V(x)=0 \quad \forall[x \in) 0,20- \\ & p_{1}(] \text { and } V(x)=0 \quad \forall[x> \\ & 20] \text { and } V(x)=0.05 \forall[x \in \\ & \left.\left(20-p_{1}, 20\right)\right] \end{aligned}$ | $p_{1} \equiv$ barrier width |


| qpl | 1D lattice with 2 frequencies/amplitudes | $\begin{array}{lll} V(x)= & p_{1} \cos \left(p_{2} x\right) & + \\ p_{3} \cos \left(p_{4} x\right) \end{array}$ | $p_{1 / 3} \equiv$ amplitudes of the lattices, $p_{2 / 4} \equiv$ frequencies of the lattices. |
| :---: | :---: | :---: | :---: |
| rot2D | 2D harmonic oscillator with rotating anisotropy | $V(x, t) \quad=$ $\frac{1}{2}((1$ + <br> $\left.p_{a}(t)\right)\left(x \cos \left(p_{1} t\right)\right.$  + <br> $\left.y \sin \left(p_{1} t\right)\right)^{2}+$ $(1$ - <br> $\left.p_{a}(t)\right)\left(y \cos \left(p_{1} t\right)\right.$  - <br> $\left.\left.x \sin \left(p_{1} t\right)\right)\right)$   | $p_{1} \equiv$ the rotation frequency, $p_{a} \equiv$ the timedependent anisotropy. $p_{2} \equiv$ anisotropy maximum, $p_{3} \equiv$ ramp-up and ramp-down time of the anisotropy, $p_{4} \equiv$ plateau time at which $p_{a}=p_{2}$. |
| stir2D | 2D harmonic trap with rotating stirring rod | $\begin{array}{ll} \hline V(x, t) & \\ \frac{1}{2}\left(x^{2}\right. & + \\ p_{3} \exp [ & \left.y^{2}\right) \\ -p_{4}\left(x-p_{2} \cos \left(p_{1} t\right)\right)^{2} & + \\ \left.\left(y-p_{2} \sin \left(p_{1} t\right)\right)^{2}\right] & \\ \hline \end{array}$ | $p_{1} \equiv$ stirring frequency, $p_{2} \equiv$ stirring radius, $p_{3} \equiv$ height of Gaussian rod, $p_{4} \equiv$ width of Gaussian rod. |
| ellipse2D | Elliptic two-dimensional hard-walled potential well | $\begin{aligned} & V(x) \\ & \begin{cases}1000 & \sqrt{\left(\frac{x}{p_{1}}\right)^{2}+\left(\frac{y}{p_{2}}\right)^{2}}>p_{3} \\ 0 & \text { else }\end{cases} \end{aligned}$ | $p_{1 / 2 / 3} \equiv$ parameters defining the ellipse. |

Table 20: Predefined potentials and parameters.

## B Predefined interaction potentials

| Which_Interaction and <br> Interaction_Type | Description | Potential |
| :---: | :---: | :---: |
| 'gauss' and 1, 2, 3, 4 | Gaussian interparticle interaction of width Interaction_Width $=\sigma$ | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0} \frac{1}{\left(\sqrt{2 \pi \sigma^{2}}\right)^{D}} \exp \left(-\frac{\|\vec{r}-\vec{r}\|^{2}}{2 \sigma^{2}}\right)$ |
| 0 | Contact interaction with constant strength $\lambda_{0}$ | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0} \delta\left(\vec{r}-\vec{r}^{\prime}\right)$ |
| 'cos' and 6 | Contact interaction with time-dependent strength $\lambda(t)$ | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0} \cos \left(I_{1} t\right) \delta\left(\vec{r}-\vec{r}^{\prime}\right)$ |
| 'sin' and 6 | Contact interaction with time-dependent strength $\lambda(t)$ | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0} \sin \left(I_{1} t\right) \delta\left(\vec{r}-\vec{r}^{\prime}\right)$ |
| 'TDHIM' and 5 | Time-dependent version of the harmonic interaction model | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0}\left(1+0.4 \sin ^{2}(t)\right)\left(\vec{r}-\vec{r}^{\prime}\right)^{2}$ |
| 'TDgauss1' and 5 | Gaussian interac- <br> tion with width <br> Interaction_Width=  <br> and time-dependent ampli-  <br> tude  | $\begin{array}{lll} W\left(\vec{r}, \vec{r}^{\prime}\right) & = & \left(\lambda_{0}\right. \\ \left.I_{1} \sin \left(I_{2} t\right)\right) \frac{1}{\left(\sqrt{2 \pi \sigma^{2}}\right)^{D}} \exp \left(-\frac{\|\vec{r}-\vec{r}\|^{2}}{2 \sigma^{2}}\right) & & \end{array}$ |
| 'TDgauss2' and 5 | Gaussian $\quad$ interac- tion with width Interaction_Width $=\sigma$ at $t=0$ that is time- dependently modulated | $\begin{aligned} & W(\vec{r}, \vec{r}) \quad=\quad \lambda_{0} \frac{1}{\left(\sqrt{2 \pi\left(\sigma^{2}+I_{1} \sin \left(I_{2} t\right)\right)}\right)^{D}} \\ & \left.\exp \left(-\frac{\|\vec{r}-\vec{r}\|^{2}}{2\left(\sigma^{2}\right.}+I_{1} \sin \left(I_{2} t\right)\right)\right) \end{aligned}$ |
| 'lennart_j' and 4 | $I_{2}$ - Screened Lennart-Jones potential | $\begin{aligned} & W\left(\vec{r}, \vec{r}^{\prime}\right)=I_{1}\left(\frac{\sigma}{\|\vec{r}-\vec{r}\|^{1}}-\frac{\sigma}{\|\vec{r}-\vec{r}\|^{6}}\right) \text { for }\|\vec{r}-\vec{r}\|>I_{2} \\ & \text { and } W\left(\vec{r}, \vec{r}^{\prime}\right)=I_{1}\left(\frac{\sigma}{I_{2}^{1} 2}-\frac{\sigma}{I_{2}^{6}}\right) \text { for }\|\vec{r}-\vec{r}\| \leq I_{2} \end{aligned}$ |
| 'HIM' and 4 | Harmonic interaction model | $W\left(\vec{r}, \vec{r}^{\prime}\right)=\lambda_{0}\left(\vec{r}-\vec{r}^{\prime}\right)^{2}$ |

Table 21: Predefined time-independent and time-dependent interaction potentials. $I_{X}$ stands for Interaction ParameterX from the input, $\sigma$ for Interaction_Width, and $D$ for the dimensionality of the problem.

## C Structure of the output of the analysis program

| Column 1 to 3 | Column 4 to $3+N_{l}$ |  |
| :--- | :--- | :--- |
| $x, y, z$ or $k_{x}, k_{y}, k_{z}$ | $\rho^{i}(x, y, z ; t)$ | or |
|  | $\rho^{i}\left(k_{x}, k_{y}, k_{z} ; t\right)$ |  |

Table 22: Structure of the <time $>N<N>M<M>x$-density. dat and $<$ time $>N<N>M<M>k$-density. dat files.
These files are generated if Density_X/Density_K it true. $x, y, z$ and $k_{x}, k_{y}, k_{z}$ are the spatial and momentum grid, respectively, $N_{l}$ is the number of internal states, and $\rho^{i}(x, y, z ; t)$ and $\rho^{i}\left(k_{x}, k_{y}, k_{z} ; t\right)$ are the spatial and momentum densities, respectively. In the case of a computation treating atoms with internal structure, the index $i$ runs through all internal states of the considered atoms and one density is output for every state. In the names of the files <time> is the time $t<\mathrm{N}\rangle$ is the particle number, <M> is the orbital number.

| Column 1 | Column 2 |  |
| :--- | :--- | :--- |
| Time $t$ | Nonescape probability |  |
|  | $P_{\text {not }}\left(t, x_{s}, x_{e}\right)$ |  |

Table 23: The nonescape probability output file Nonescape.
If the input variable Pnot and the borders $x_{s}$ and $x_{e}$ were defined with xstart and xend in the input of the analysis program, this file is created.

| Column 1 | Column 2 |  |  | Column 3 | Column 4 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Time $t$ | $\begin{aligned} & S_{\rho-r}(t) \\ & -\int d \vec{r} \rho(\vec{r} ; t) \ln [\rho(\vec{r} ; t)] \end{aligned}$ |  |  | $\begin{aligned} & S_{\rho-k}(t) \\ & -\int d \vec{k} \rho(\vec{k} ; t) \ln [\rho(\vec{k} ; t)] \end{aligned}$ | $\begin{array}{ll} S_{C}(t) & = \\ \sum_{\vec{n}}-\left\|C_{\vec{n}}(t)\right\|^{2} \ln \left[\left\|C_{\vec{n}}(t)\right\|^{2}\right] \end{array}$ |
| Column 5 |  | Column 6 | Column 7 |  |  |
| $\begin{aligned} & S_{n}(t) \\ & \sum_{i}-\frac{n_{i}(t)}{N} l n \end{aligned}$ |  | $\begin{aligned} & I \quad= \\ & \frac{\sum_{\vec{n}}\left\|C_{\vec{n}}(t)\right\|^{4}}{}= \end{aligned}$ |  | $-\left\|C_{\vec{n}}(t)\right\|^{2} \ln \left[\left\|C_{\vec{n}^{\prime}}(t)\right\|^{2}\right]$ |  |

Table 24: Structure of the Entropy.dat file.
This file is generated when Entropy is set to true. The last column is only present if NBody_C_Entropy=.T. is set in analysis.inp.

| Column 1 | Column 2 | Column 3 | Column 4 |
| :--- | :--- | :--- | :--- |
| Time $t$ | $S_{\rho^{(2)}-r}(t)$ | $S_{\rho-k}(t)$ | $S_{\rho_{k}^{(G O)}(t)}$ |
|  | $-\int d \vec{r}_{1} d \vec{r}_{2} \rho^{(2)}\left(\vec{r}_{1}, \vec{r}_{2} ; t\right)$ | $-\int d \vec{k}_{1} d \vec{k}_{2} \rho^{(2)}\left(\vec{k}_{1}, \vec{k}_{2} ; t\right)$ | $\sum_{i}-\frac{\rho_{i}^{(G O)}(t)}{N} \ln \left[\frac{\rho_{i}^{(G O)}(t)}{N}\right]$ |
|  | $\ln \left[\rho^{(2)}\left(\vec{r}_{1}, \vec{r}_{2} ; t\right)\right]$ | $\ln \left[\rho^{(2)}\left(\vec{k}_{1}, \vec{k}_{2} ; t\right)\right]$ |  |

Table 25: Structure of the TwoBody_Entropy . dat file
. This file is generated when TwoBody_Entropy is true.

| Column 1 to 3 | $\begin{aligned} & \text { Column } 4 \\ & \text { to } 6 \end{aligned}$ | $\begin{aligned} & \text { Column } 7+ \\ & 5(i-1) \\ & \hline \end{aligned}$ | $\begin{aligned} & \text { Column } 8+5(i-1) \& \\ & 9+5(i-1) \end{aligned}$ | $\begin{array}{ll} \text { Column } 10+ \\ 5(i-1) & \\ \hline \end{array}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & x, y, z \text { or } \\ & k_{x}, k_{y}, k_{z} \end{aligned}$ | $\begin{aligned} & x^{\prime}, y^{\prime}, z^{\prime} \text { or } \\ & k_{x}^{\prime}, k_{y}^{\prime}, k_{z}^{\prime} \end{aligned}$ | $\rho_{i}(x, y, z ; t)$ <br> or $\rho_{i}\left(k_{x}, k_{y}, k_{z} ; t\right)$ | $\begin{aligned} & \rho_{i}^{(1)}\left(x, y, z \mid x^{\prime}, y^{\prime}, z^{\prime} ; t\right) \text { or } \\ & \rho_{i}^{(1)}\left(k_{x}, k_{y}, k_{z} \mid k_{x}^{\prime}, k_{y}^{\prime}, k_{z}^{\prime} ; t\right) \end{aligned}$ | $\rho_{i}\left(x^{\prime}, y^{\prime}, z^{\prime} ; t\right)$ <br> or $\rho_{i}\left(k_{x}^{\prime}, k_{y}^{\prime}, k_{z}^{\prime} ; t\right)$ |
| $\begin{aligned} & \rho_{i}^{(2)}\left(x, y, z \mid x^{\prime}, y^{\prime}, z^{\prime} ; t\right) \text { or } \\ & \rho_{i}^{(2)}\left(k_{x}, k_{y}, k_{z} \mid k_{x}^{\prime}, k_{y}^{\prime}, k_{z}^{\prime} ; t\right) \end{aligned}$ |  |  |  |  |

Table 26: Structure of the <time $>N<N>M<M>x$-correlations.dat and <time $>N<N>M<M>k$-correlations.dat files for multilevel computations.
(For the explanation of the filenames, see table 22). These files are created, if the input variable Correlations_X and Correlations_K, respectively, are set to be true. The files contain all necessary quantities to compute the one-body as well as the diagonal of the two-body normalized (Glauber-) correlation function $g_{i}^{(1)}$ and $g_{i}^{(2)}$, respectively, for all internal states $i$. For instance, $\left|g_{1}^{(1)}\right|^{2}=\left|\frac{\rho_{1}^{(1)}\left(x_{1}, x_{1}^{\prime} ; t\right)}{\sqrt{\rho_{1}\left(x_{1} ; t\right) \rho_{1}\left(x_{1}^{\prime} ; t\right)}}\right|^{2}$ can be plotted as the value of $\left.(\text { (Column } 8)^{2}+(\text { Column } 9)^{2}\right)$ divided by (Column 7) $\times$ (Column 10).


Table 27: Structure of the <time $>N<N>M<M>x$-correlations.dat and <time $>\mathrm{N}<\mathrm{N}>\mathrm{M}<\mathrm{M}>\mathrm{k}$-correlations.dat files.
(For the explanation of the filenames, see table 22). These files are created, if the input variable Correlations_X and Correlations_K, respectively, are set to be true. The files contain all necessary quantities to compute the one-body as well as the diagonal of the two-body normalized (Glauber-) correlation function $g^{(1)}$ and $g^{(2)}$, respectively. For instance, $\left|g^{(1)}\right|^{2}=\left|\frac{\rho^{(1)}\left(x_{1}, x_{1}^{\prime} ; t\right)}{\sqrt{\rho\left(x_{1} ; t\right) \rho\left(x_{1}^{\prime} ; t\right)}}\right|^{2}$ can be plotted as the value of $\left((\operatorname{Column} 8)^{2}+(\operatorname{Column} 9)^{2}\right)$ divided by (Column 7$) \times($ Column 10).

| Column $1 \& 2$ | Column $3 \& 4$ | Column 5 | Column 6 |
| :--- | :--- | :--- | :--- |
| $x, x^{\prime}$ or $k, k^{\prime}$ | $\rho^{(1 / 2)}\left(x \mid x^{\prime} ; t\right)$ or $\rho^{(1 / 2)}\left(k \mid k^{\prime} ; t\right)$ | $\rho(x ; t)$ or $\rho(k ; t)$ | $\rho\left(x^{\prime} ; t\right)$ or $\rho\left(k^{\prime} ; t\right)$ |

Table 28: Structure of the <time $>N<N>M<M><x / k>c o r r<1 / 2>$ restr . dat files.
The generation of these files is triggered by the analysis input variables corr1restr,corr2restr,corr1restrmom,corr2restrmom. Similar to the above table 27, the normalized correlation functions $g^{(1)}$ and $g^{(2)}$ can be computed from the contents of these files, but for one-dimensional computations and on a restricted grid which is specified through the analysis input variables <x/k>ini<1/2>, <x/k>fin<1/2>, <x/k>pts<1/2>, respectively.

| Column 1 to 3 | Column 4 | Column 5 | Column 6\& 7 |
| :---: | :---: | :---: | :---: |
| $\begin{aligned} & x, y, z \text { or } \\ & k_{x}, k_{y}, k_{z} \end{aligned}$ | $\begin{aligned} & \hline \rho\left(x_{r e f}, y_{r e f},\right. \\ & \text { or } \\ & \rho\left(k_{x, r e f}, k_{y, r}\right. \end{aligned}$ | $\begin{aligned} & \text { zq }(\not(\not \approx ; \text {,ty }), z ; t) \text { or } \\ & \rho\left(k_{x}, k_{y}, k_{z} ; t\right) \\ & \text { ref } \left., k_{z, r e f} ; t\right) \end{aligned}$ | $\left.\begin{array}{l} \rho^{\text {order }}\left(\left\{k_{x, \text { ref }}, k_{y, \text { ref }}, k_{z, \text { ref }}\right\}\right. \\ \text { or } \\ \rho^{\text {order }}\left(\left\{k_{x, \text { ref }}, k_{y, \text { ref }}, k_{z, \text { ref }}\right\}\right. \end{array}\right\}$ |

Table 29: Structure of the <time>N<N>M<M>x/k-order-<order>-correlations1D.dat (For the explanation of the filenames, see table 22). These files are created, if the input variable anyordercorrelations_X and anyordercorrelations_K as well as oneD, respectively, are set to be true. The files contain all necessary quantities to compute the correlation functions up to order <order>. Here, $x_{r e f} / y_{r e f} / z_{r e f} / k_{x, r e f} / k_{y, r e f} / k_{z, \text { ref }}$ correspond to the reference point specified in the input file by c_ref_x/c_ref_y/c_ref_z, respectively.

| Column 1 | Column 2 \& 3 |
| :--- | :--- |
| time $t$ | Real \& imaginary part of $\tau=\frac{\left\langle x_{1} x_{2}\right\rangle-\langle x\rangle^{2}}{\left\langle x^{2}\right\rangle-\langle x\rangle^{2}}$ |

Table 30: Structure of the CorrelationCoefficient.dat file. This file is created when Correlation_Coefficient it true.

| Column 1 to 3 | Column 4 and 5 | Column 6 and 7 | Column 8 to 10 | Column 11 \& 12 |
| :---: | :---: | :---: | :---: | :---: |
| $x, y, z$ | Real and imaginary part of $\rho^{(2)}\left(\vec{r}_{1}=\vec{r}_{1}^{\prime}=\right.$ $\vec{R}, \vec{r}_{2}=\vec{r}_{2}^{\prime}=\vec{r}$ ) | Real and imaginary part of $\rho^{(1)}\left(\vec{r}_{1}=\vec{R}, \vec{r}_{1}^{\prime}=\right.$ $\vec{r}$ ) | $k_{x}, k_{y}, k_{z}$ | Real and imaginary part of dynamic structure factor $G=1+N \mathcal{F}$ $\left(\rho^{(2)}\left(\vec{r}_{1}=\vec{r}_{1}^{\prime}=\vec{R}\right.\right.$, $\left.\overrightarrow{r_{2}}=\vec{r}_{2}^{\prime}=\vec{r}-1\right)$ |

Table 31: Structure of the <time $>N<N>M<M>x$-StructureFactor.dat files. These files are output if StructureFactor is true.

| Column 1 | Column 2 | Column 3 | Column 4 |
| :--- | :--- | :--- | :--- |
| Time $t$ | $P_{0}^{2}(t)$ | $P_{1}^{1}(t)$ | $P_{2}^{0}(t)$ |

Table 32: Structure of the lossops_N2_<border>. dat files.
The generation of such a file is triggered by the lossops input variable being set to .T.. <border> is controlled by the border input variable. For each point in time $t$, a line in this file contains the probability $P_{0}^{2}(t)$ to find 2 particles to the left of border, the probability $P_{1}^{1}(t)$ to find one particle to the left and one to the right of border, and the probability $P_{2}^{0}(t)$ to find two particles to the right of border.

| Column 1 to 4 | Column 5 | Column 6 | Column 7 \& 8 | Column 9 |
| :--- | :--- | :--- | :--- | :--- |
| $r_{1 x}, r_{1 y}, r_{2 x}, r_{2 y}$ | $\rho\left(\vec{r}_{1} ; t\right)$ | $\rho\left(\vec{r}_{2} ; t\right)$ | $\rho^{(1)}\left(\vec{r}_{1} \mid \vec{r}_{2} ; t\right)$ | $\rho^{(2)}\left(\vec{r}_{1} \mid \vec{r}_{2} ; t\right)$ |

Table 33: The structure of the <time>N<N>M<M><x/k><Slice 1>-<Slice 2>-correlations.dat files
. These are output by the analysis program if the analysis input variable MOMSPACE2D or REALSPACE2D is set to .T.. <Slice $1 / 2>$ specify which cut through the real- or momentum-space density are in the file.

| Column 1 and 2 | Column 3 | Column 4 and 5 | Column 6 | Column 7 |
| :--- | :--- | :--- | :--- | :--- |
| $r_{1 x}, r_{1 y}$ | $\rho\left(\vec{r}_{1} ; t\right)$ | $\rho^{(1)}\left(\vec{r}_{1} \mid-\vec{r}_{1} ; t\right)$ | $\rho\left(-\vec{r}_{1} ; t\right)$ | $\rho^{(2)}\left(\vec{r}_{1},-\vec{r}_{1} ; t\right)$ |

Table 34: The structure of the <time>N<N>M<M><x/k>-SkewCorrelations.dat files
These files are output of the analysis program if the analysis input variable MOMSKEW2D or REALSKEW2D is set to .T. .

| Column $1,2,3$ | Column 4 to $3+N_{\text {shots }}$ |
| :--- | :--- |
| $x, y, z$ | $N_{\text {shots s samples of the }} N$-body <br> density |

Table 35: The structure of the <time>N<N>M<M><x/k>SingleShots.dat files . These files are output of the analysis program if the analysis input variable SingleShot_Analysis or SingleShot_FTAnalysis is set to .T. .

| Column $1,2,3$ | Column 4 |
| :--- | :--- |
| $x, y, z$ | Histogram of $N_{\text {samples }}$ samples of <br> the centre-of-mass-operator. |

Table 36: The structure of the <time>N<N>M<M><x/k>CentreOfMass.dat files
. These files are output of the analysis program if the analysis input variable CentreOfMass or CentreOfMomentum is set to .T. .

| Column 1 <br> to 3 | Column 4 | Column 5 <br> to $(5+M)$ | Column $(6+M) \&$ <br> $(7+M)$ |  <br> $(9+M)$ to $(8+3 M)$ <br>  <br> $\&(9+3 M)$ |
| :--- | :--- | :--- | :--- | :--- |
| $x, y, z$ | $\xi_{\text {avg }}(x, y, z ; t)$ | $\xi_{1}(x, y, z ; t)$ <br> to | $\nabla_{x} \xi_{\text {avg }}(x, y, z ; t)$ <br> $\nabla_{y} \xi_{\text {avg }}(x, y, z ; t)$ | $\nabla_{x} \xi_{1}(x, y, z ; t)$ $\&$ <br> $\nabla_{y} \xi_{1}(x, y, z ; t)$ to <br> $\nabla_{x} \xi_{M}(x, y, z ; t)$ $\&$ <br> $\nabla_{y} \xi_{M}(x, y, z ; t)$  |

Table 37: The structure of the <time $>\mathrm{N}<\mathrm{N}>\mathrm{M}<\mathrm{M}>$ phase. dat files.
The generation of the files is toggled by setting the analysis input variable PHASE to .T.. If additionally GRADIENT it set to .T. then, Columns $(8+M)$ to $(9+3 M)$ containing the phase gradients will be generated. Here $x, y, z$ are the coordinates, $\xi_{\text {avg }}(x, y, z ; t)$ is the average phase, $\xi_{1}(x, y, z ; t)$ to $\xi_{M}(x, y, z ; t)$ are the $M$ orbital phases, $\nabla_{x} \xi_{\text {avg }}(x, y, z ; t) \& \nabla_{y} \xi_{\text {avg }}(x, y, z ; t)$ is the $x$ and $y$ component of the average phase's gradient, and $\nabla_{x} \xi_{1}(x, y, z ; t) \& \nabla_{y} \xi_{1}(x, y, z ; t)$ to $\nabla_{x} \xi_{M}(x, y, z ; t) \&$ $\nabla_{y} \xi_{M}(x, y, z ; t)$ are the $x$ and $y$ components of the gradients of the $M$ orbital phases.


[^0]:    ${ }^{1}$ Fast Fourier Transform
    ${ }^{2}$ Discrete Variable Representation

